



BECQUEREL
PROJECT

Проект
БЕККЕРЕЛЬ

Beryllium (Boron)

Clustering

Quest in

Relativistic Multifragmentation

<http://becquerel.lhe.jinr.ru>

Light nucleus clustering in fragmentation above 1 A GeV

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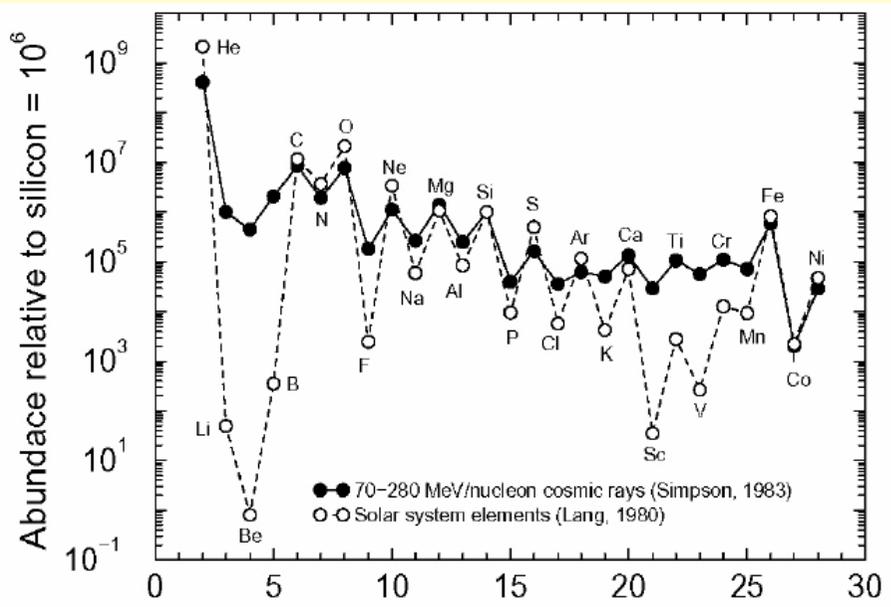
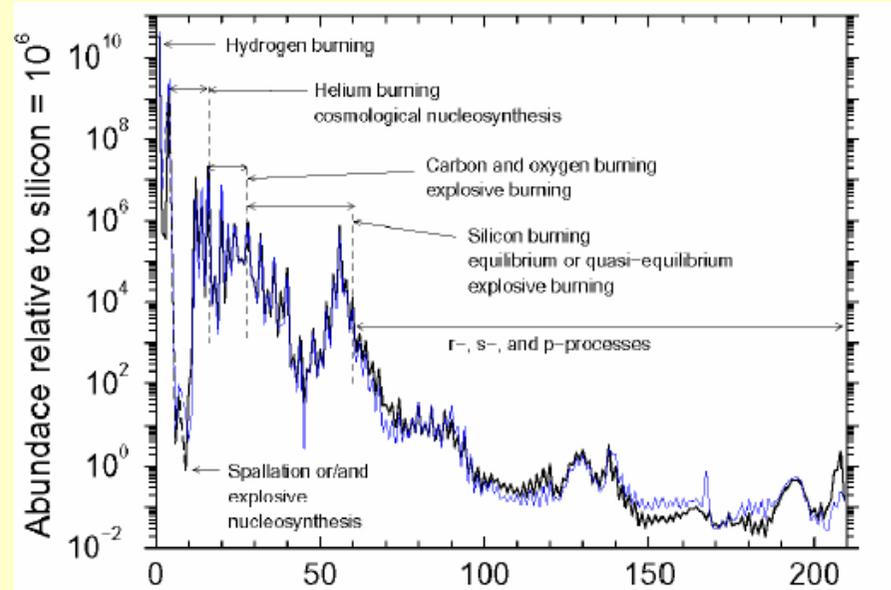
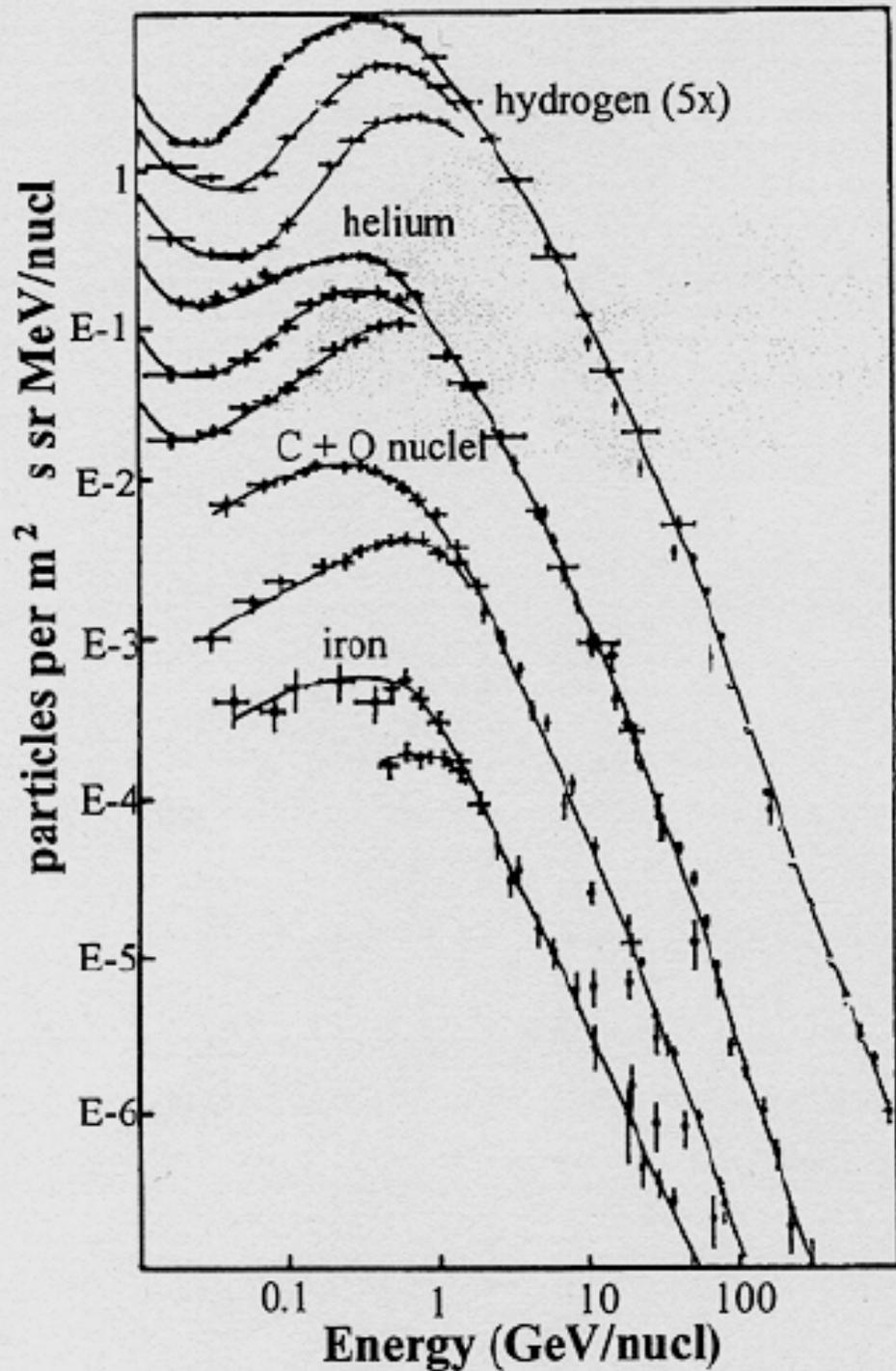
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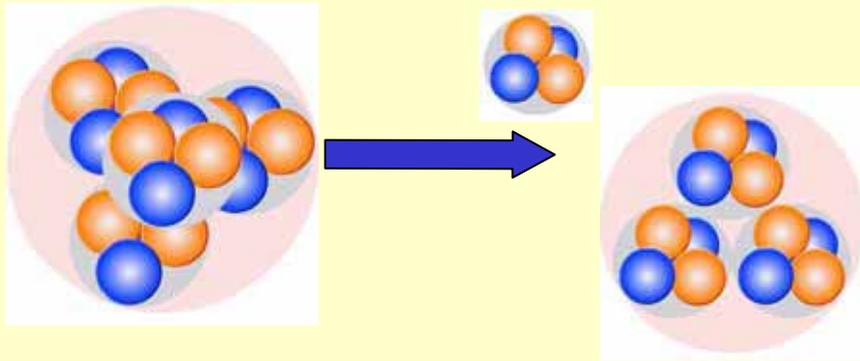
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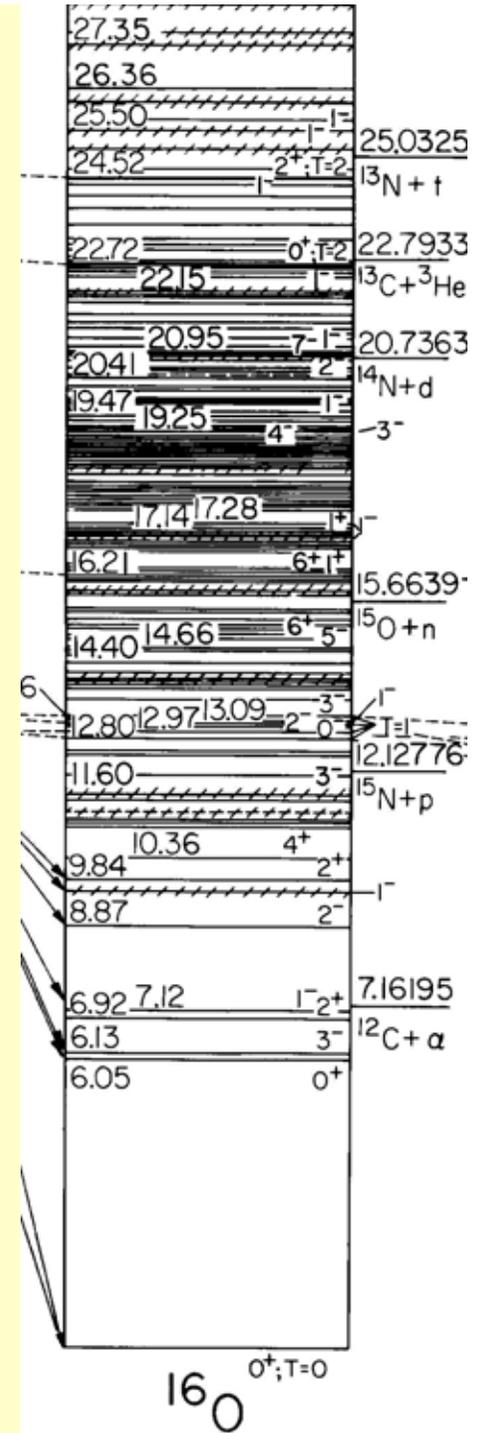
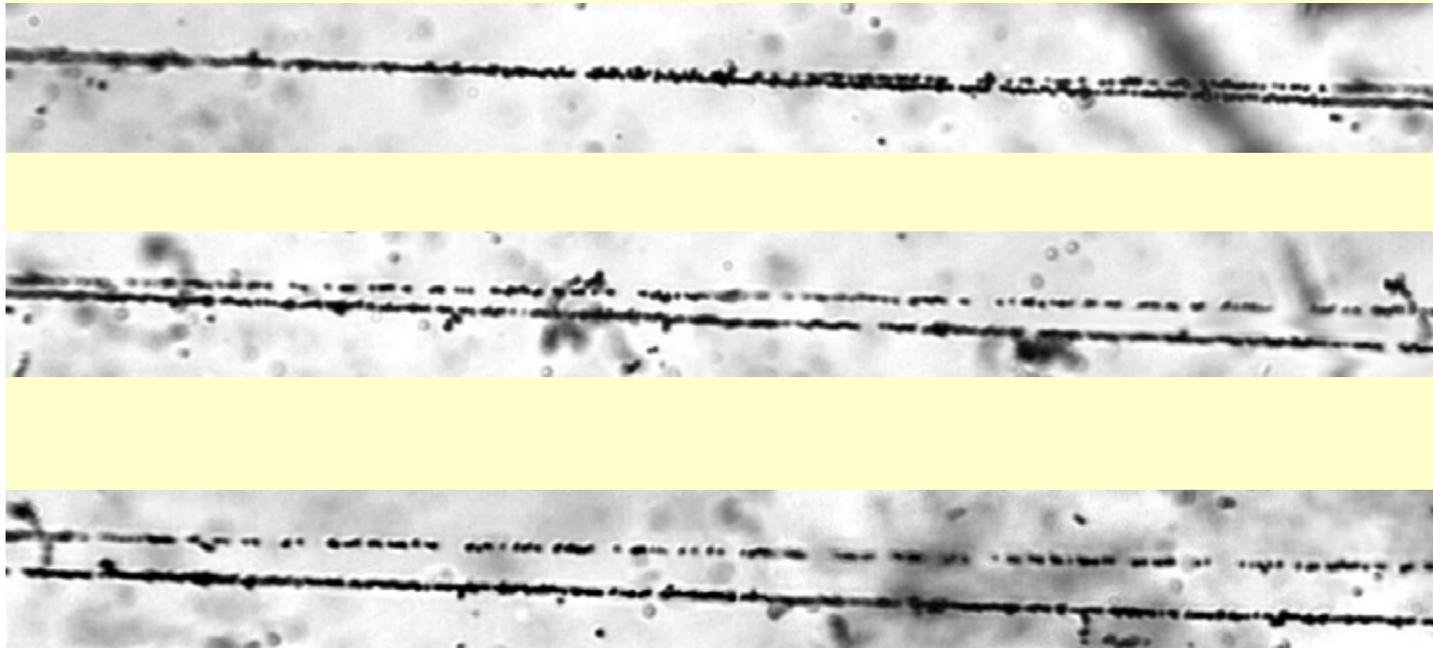
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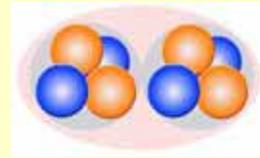
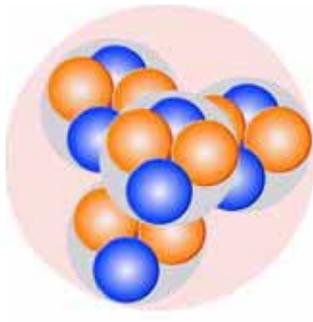
^h*P. J. Safarik University, Kosice, Slovakia*



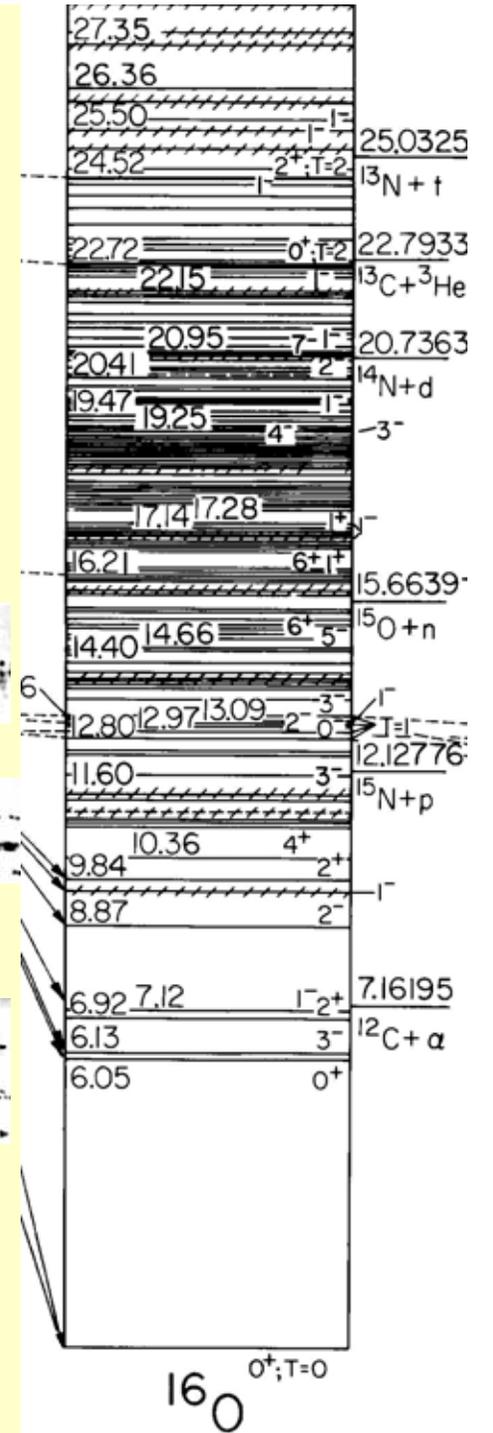
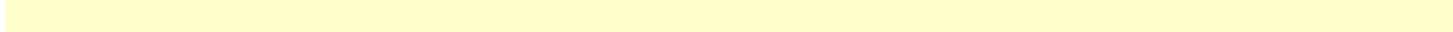
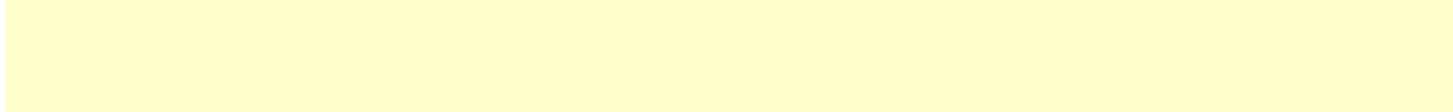
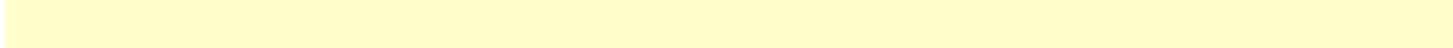
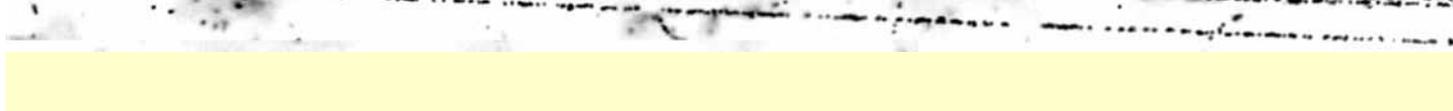
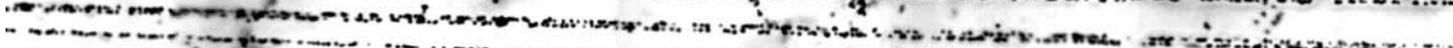
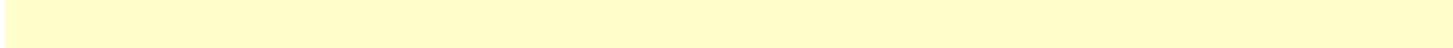
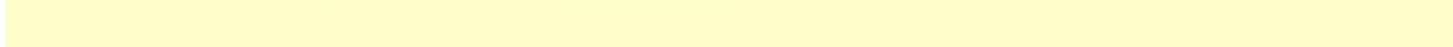
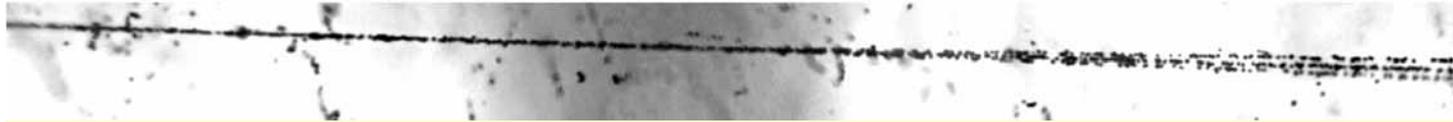


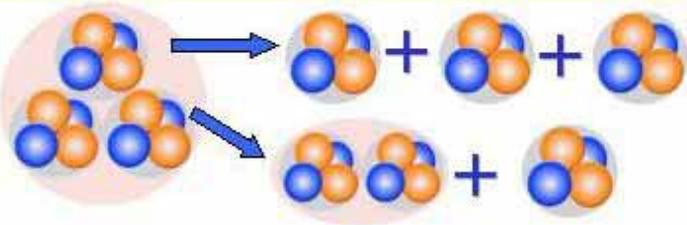
4.5 A GeV/c ^{16}O



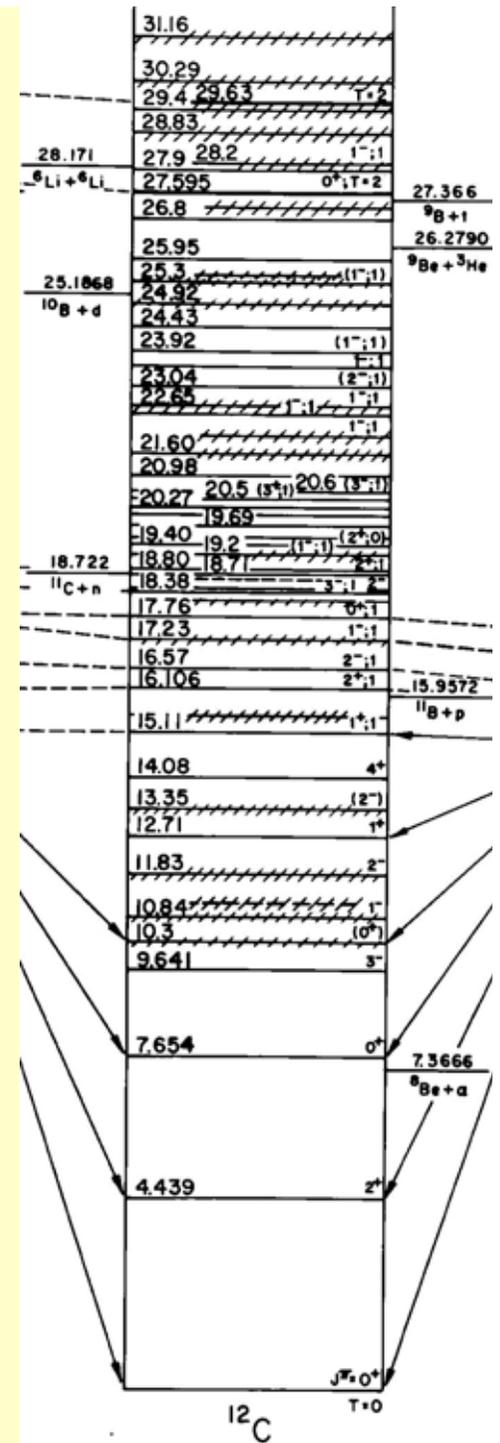


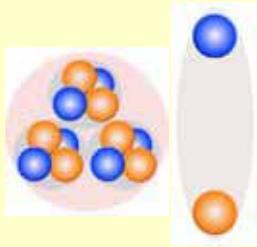
4.5A GeV/c ¹⁶O



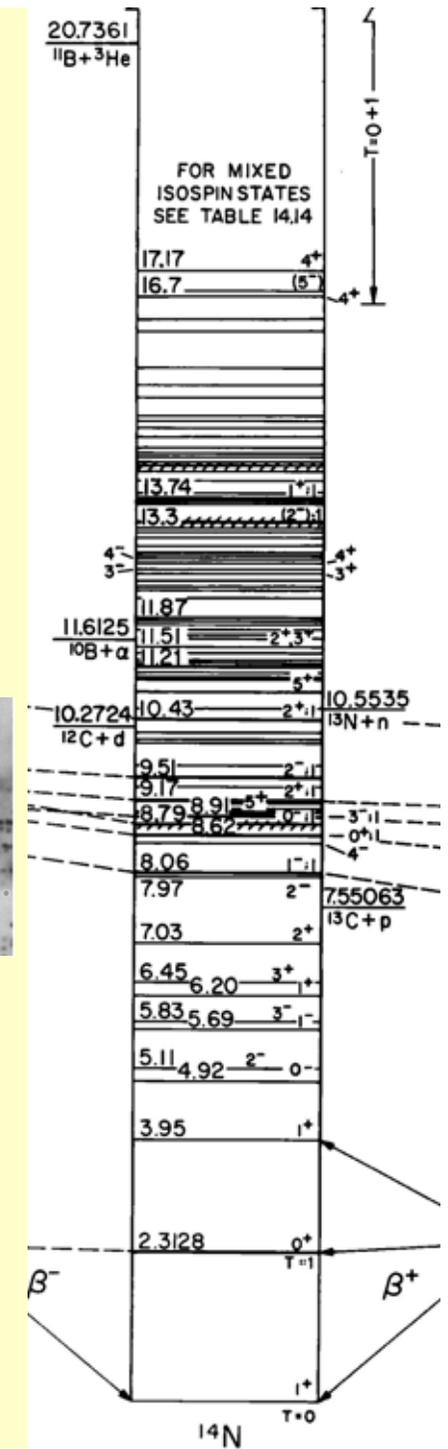
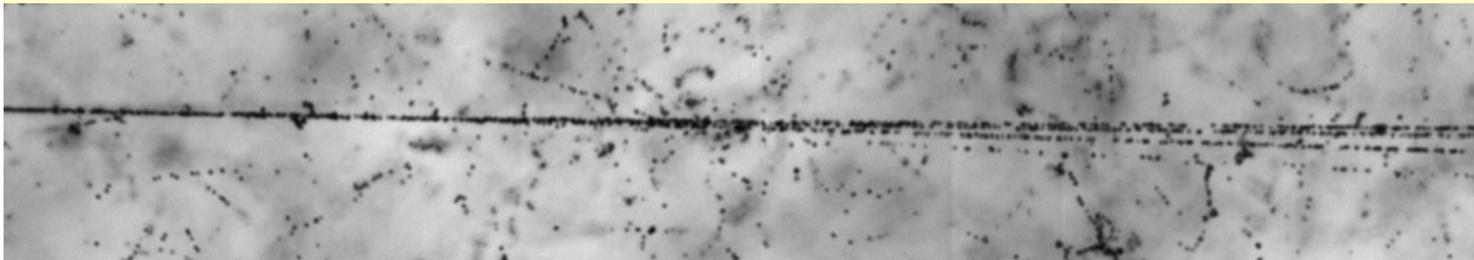


**4.5A GeV/c ^{12}C
Coherent Dissociation**





2.9A GeV/c ^{14}N



PROGRESS
in
COSMIC RAY PHYSICS

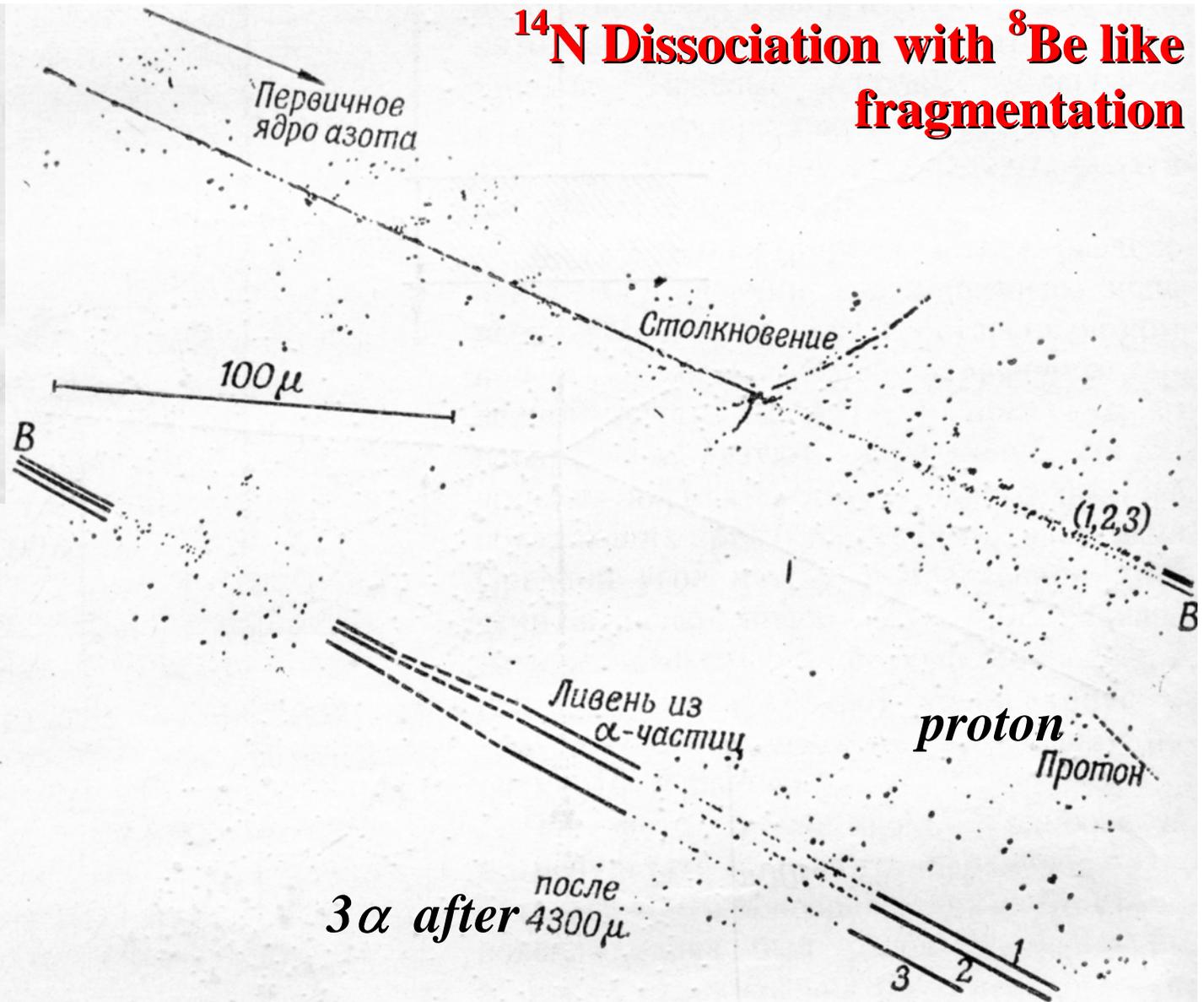
Edited by
J. G. WILSON

Contributors

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| W. O. Lock | B. Peters | N. Dallaporta |
| D. N. Perkins | H. V. Neher | E. P. George |
| C. C. Butler | | H. Elliot |

AMSTERDAM, 1952

^{14}N Dissociation with ^8Be like fragmentation



Фиг. 7. Ядро азота столкнулось с ядром эмульсии. Повидимому, произошло скользящее столкновение, при котором заряд первичного ядра уменьшается на единицу. Остаток, представляющий собой возбужденное ядро углерода, распадается затем на 3 α -частицы, которые в лабораторной системе испускаются в узком конусе в направлении движения первичной частицы

PROGRESS
in
COSMIC RAY PHYSICS

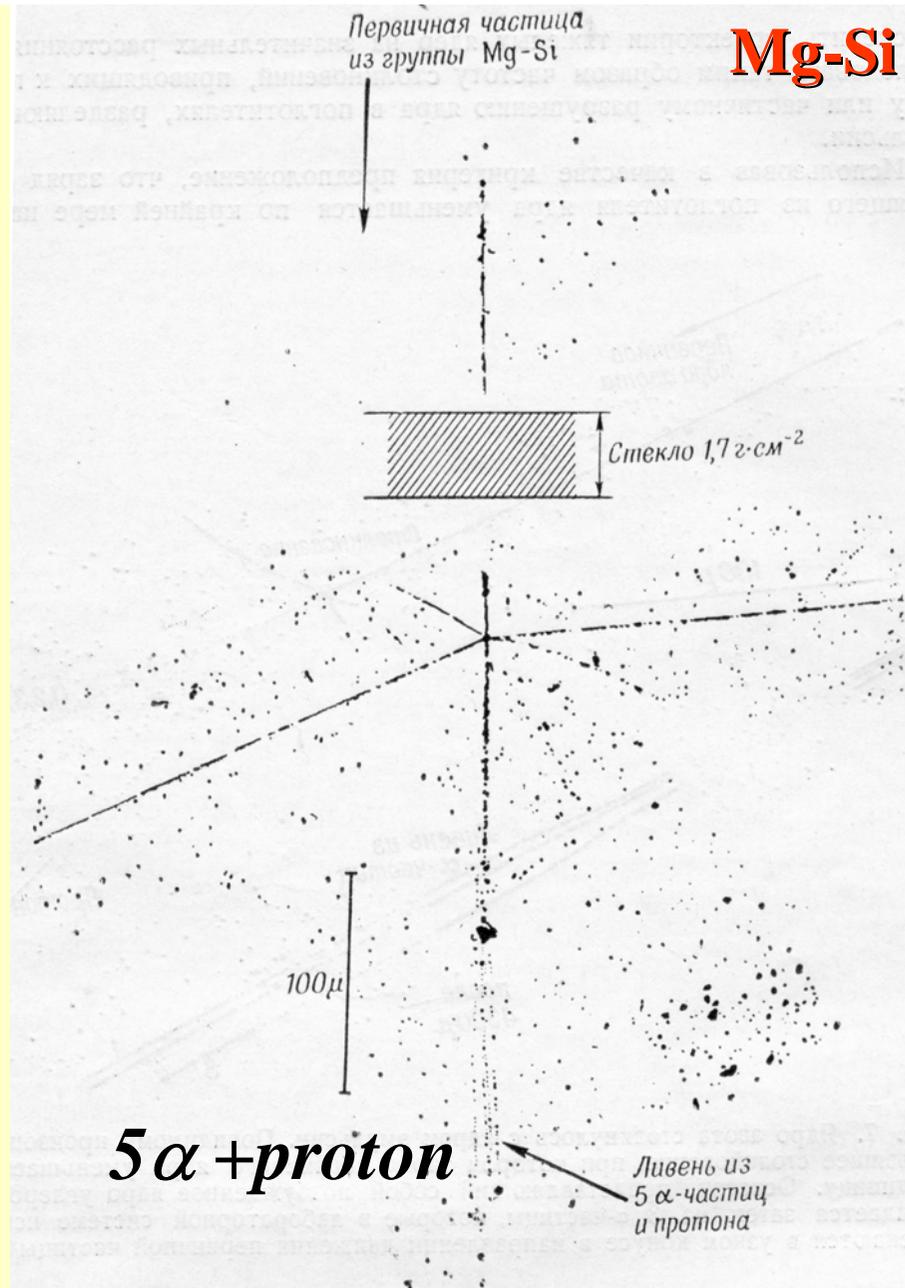
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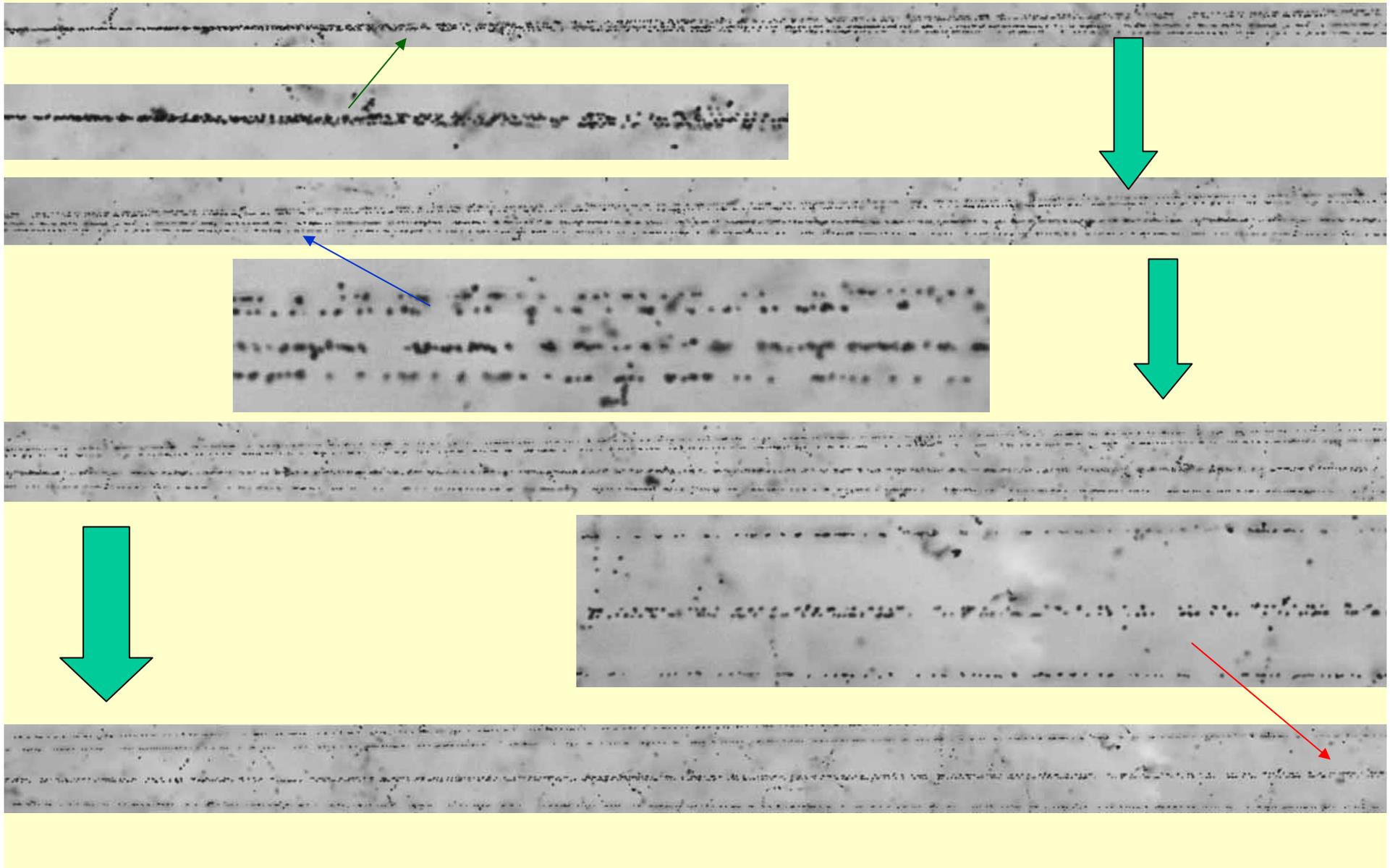
**Mg-Si Dissociation into
charge state
2+2+2+2+2+1**



Фиг. 6. Ядро из группы Mg—Si столкнулось с ядром эмульсии. Предполагают, что узкий ливень, состоящий из протона и 5 α -частиц, возник в результате испарения первичного ядра, возбужденного столкновением. Сстальные частицы, испущенные в звезде, являются, видимо, осколками ядра мишени.

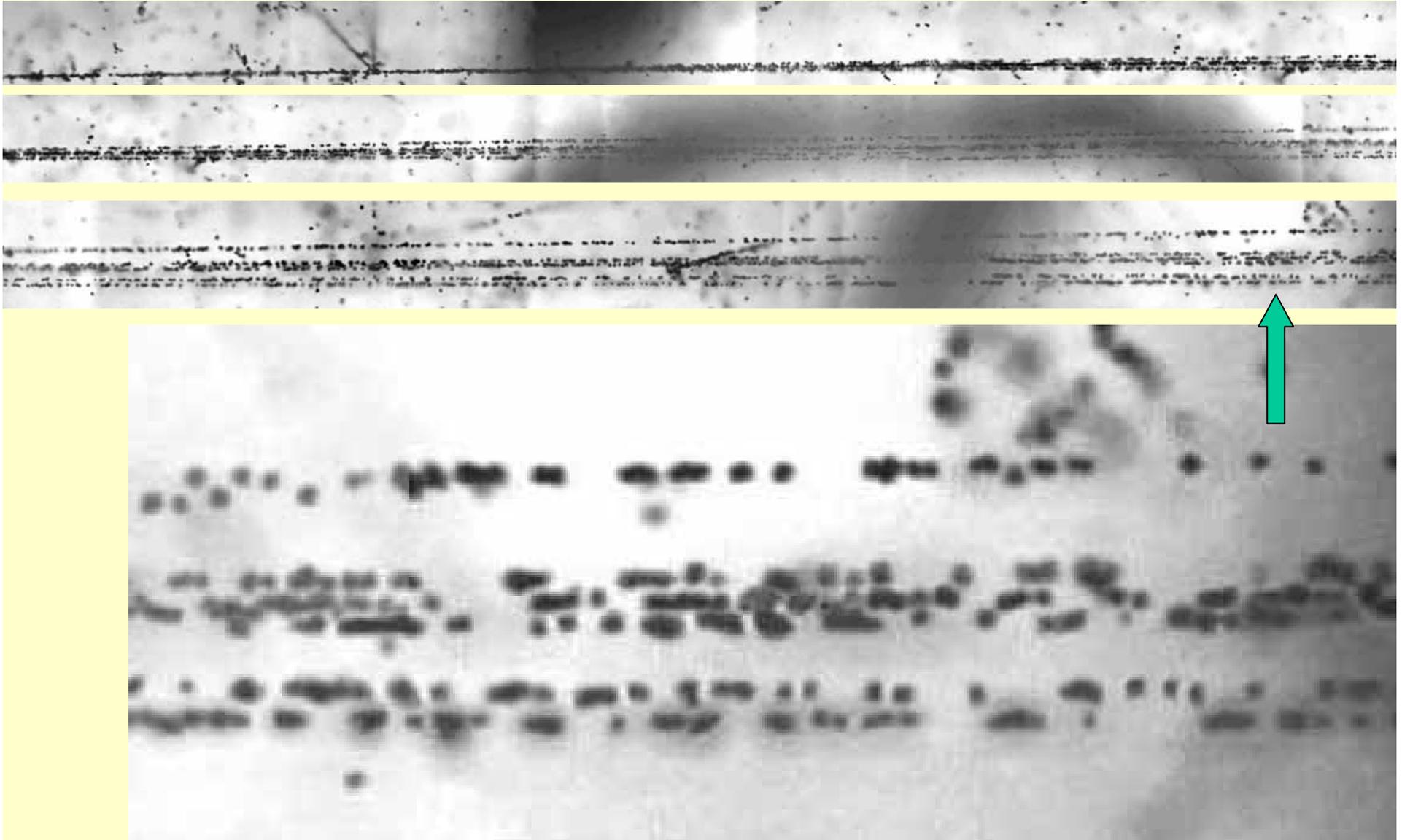
4.5A GeV/c ^{20}Ne Peripheral Dissociation into charge state

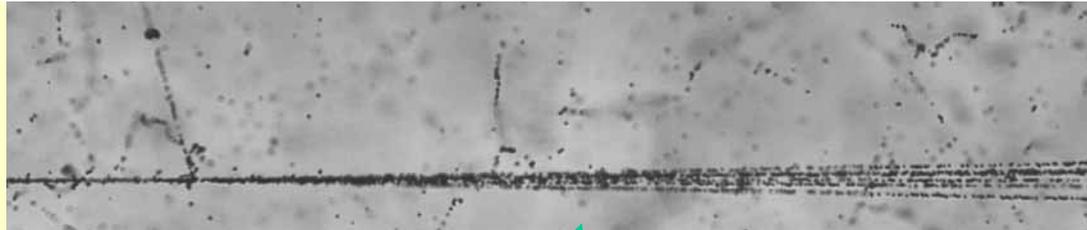
2+2+2+2+2 with ^8Be like fragments



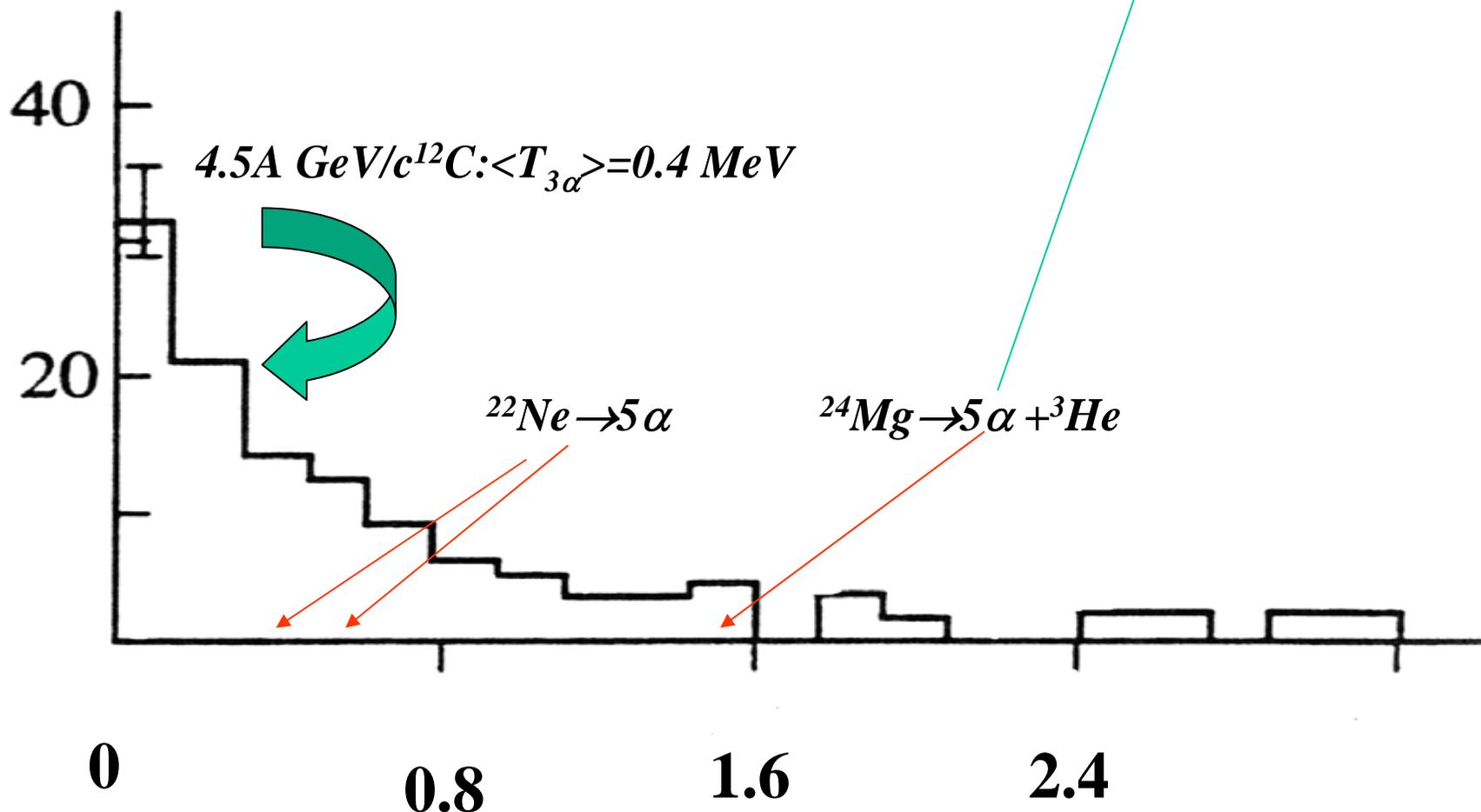
4.5A GeV/c ^{24}Mg Peripheral Dissociation into charge state

2+2+2+2+2+2 with ^8Be and $^{12}\text{C}^*$ like fragments





$dN/dT_{n\alpha}$



$$T_{n\alpha} = (M_{n\alpha}^* - n_{\alpha} M_{\alpha}) / (4 n_{\alpha}), \text{ MeV}$$

Boltzmann constant, k /approx 10^{-4} eV K⁻¹
Typical Temperature Range, T /approx $5 \cdot 10^{8-9}$ K per α

$$p_\alpha = \sqrt{2m_\alpha \cdot T_\alpha} \quad p_\alpha \text{ /approx } 20\text{-}120 \text{ MeV}$$

Planck constant, \hbar /approx 200 MeV fm

$\lambda = \hbar/p$ de Broglie wave lengths /approx 1-10 fm

$$\lambda_\alpha^{\text{coh}} \text{ /approx } R_\alpha \quad \lambda_{\text{He}}^{\text{coh}} \text{ /approx } R_{\text{He}}$$

$$T_\alpha/T_{\text{He}} = T_\alpha/T_{\text{He}} = (R_{\text{He}}/R_\alpha)^2 \text{ /approx } 10^{10}$$

Macroscopic quantum coherence phenomena in atomic physics /approx 1 K

Macroscopic quantum coherence phenomena in atomic physics /approx 10^{10} K

Topology of “White Stars” in the Relativistic Fragmentation of Light Nuclei

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Abstract—Experimental observations of the multifragmentation of relativistic light nuclei by means of emulsions are surveyed. Events that belong to the type of “white stars” and in which the dissociation of relativistic nuclei is not accompanied by the production of mesons and target-nucleus fragments are considered. An almost complete suppression of the binary splitting of nuclei to fragments of charge in excess of two, $Z > 2$, is a feature peculiar to charge topology in the dissociation of Ne, Mg, Si, and S nuclei. An increase in the degree of nuclear fragmentation manifests itself in the growth of the multiplicity of singly and doubly charged fragments ($Z = 1, 2$) as the charge of the unexcited fragmenting-nucleus part (which is the main part) decreases. Features of the production of systems formed by extremely light nuclei α , d , and t are studied in the dissociation of the stable isotopes of Li, Be, B, C, N, and O to charged fragments. Manifestations of ^3He clustering can be observed in “white stars” in the dissociation of neutron-deficient isotopes of Be, B, C, and N. © 2005 Pleiades Publishing, Inc.

Table 1. Charge-topology distribution of white stars in the dissociation of ^{24}Mg nuclei with an energy of 3.65 GeV per nucleon

| | | | | | | | | | | | | | | | | | | | | |
|-----------|----|----|----|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|
| Z_f | 11 | 10 | 10 | 9 | 9 | 8 | 8 | 8 | 7 | 7 | 6 | 5 | 5 | 5 | 4 | 4 | 3 | - | - | - |
| $N_{Z=1}$ | 1 | 2 | - | 3 | 1 | 4 | 2 | - | 3 | 1 | 2 | 5 | 3 | 1 | 6 | 4 | 5 | 6 | 4 | 2 |
| $N_{Z=2}$ | - | - | 1 | - | 1 | - | 1 | 2 | 1 | 2 | 2 | 1 | 2 | 3 | 1 | 2 | 2 | 3 | 4 | 5 |
| N_{ev} | 10 | 14 | 8 | 5 | 9 | 1 | 7 | 4 | 4 | 2 | 4 | 2 | 1 | 1 | 2 | 1 | 3 | 1 | 2 | 2 |

Table 2. Charge-topology distribution of white stars in the dissociation of ^{22}Ne nuclei with an energy of 3.3 GeV per nucleon

| | | | | | | | | | | | | | |
|-----------|----|----|---|---|---|---|---|---|-----|---|-----|---|---|
| Z_f | 9 | 8 | 8 | 7 | 6 | 6 | 5 | 5 | 5+3 | 4 | 4+3 | - | - |
| $N_{Z=1}$ | 1 | - | 2 | 1 | - | 2 | 1 | 3 | - | - | 3 | 2 | - |
| $N_{Z=2}$ | - | 1 | - | 1 | 2 | 1 | 2 | 1 | 1 | 3 | - | 4 | 5 |
| N_{ev} | 22 | 51 | 6 | 7 | 5 | 2 | 1 | 1 | 1 | 2 | 1 | 1 | 3 |

Table 3. Charge-topology distribution of white stars in the dissociation of ^{28}Si nuclei with an energy of 3.65 GeV per nucleon

| | | | | | | | | | | | | | | | | | | | | | | | | | | | |
|-----------|----|----|----|----|----|----|----|----|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|---|----|
| Z_f | 13 | 12 | 12 | 11 | 11 | 10 | 10 | 10 | 9 | 9 | 9 | 8 | 8 | 8 | 7 | 7 | 7 | 6 | 6 | 6 | 6 | 5 | 5 | 4 | - | - | - |
| $N_{Z=1}$ | 1 | - | 2 | 1 | 3 | - | 2 | 4 | 1 | 3 | 5 | 6 | 2 | 4 | 3 | 5 | 7 | 2 | 4 | 6 | 8 | 3 | 5 | 2 | 2 | 8 | 10 |
| $N_{Z=2}$ | - | 1 | - | 1 | - | 2 | 1 | - | 2 | 1 | - | - | 2 | 1 | 2 | 1 | - | 3 | 2 | 1 | - | 3 | 2 | 4 | 6 | 3 | 2 |
| N_{ev} | 9 | 3 | 15 | 11 | 6 | 2 | 7 | 2 | 2 | 8 | 3 | 2 | 5 | 6 | 1 | 3 | 3 | 3 | 5 | 8 | 1 | 1 | 3 | 1 | 1 | 2 | 3 |

Table 4. Charge-topology distribution of white stars in the dissociation of ^{32}S nuclei with an energy of 200 GeV per nucleon

| | | | | | | | | | | | | | | | | | | |
|-----------|----|----|----|----|----|----|----|----|----|----|----|----|---|---|---|-----|---|-----|
| Z_f | 15 | 14 | 14 | 13 | 13 | 12 | 12 | 11 | 11 | 10 | 10 | 10 | 9 | 8 | 8 | 7+3 | 7 | 5+3 |
| $N_{Z=1}$ | 1 | - | 2 | 1 | 3 | 2 | 4 | 3 | 5 | 2 | 4 | 6 | 3 | - | 6 | 4 | 3 | 4 |
| $N_{Z=2}$ | - | 1 | - | 1 | - | 1 | - | 1 | - | 2 | 1 | - | 2 | 4 | 1 | 1 | 3 | 2 |
| N_{ev} | 99 | 11 | 48 | 7 | 6 | 3 | 4 | 4 | 1 | 1 | 2 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |

Table 5. Charge-topology distribution of white stars in the dissociation of ^{16}O nuclei with an energy of 3.65 GeV per nucleon

| | | | | | | | | | |
|-----------|----|---|----|---|----|---|---|---|---|
| Z_f | 7 | 6 | 6 | 5 | 5 | 4 | 4 | – | – |
| $N_{Z=1}$ | 1 | 2 | – | 3 | 1 | – | 2 | – | 2 |
| $N_{Z=2}$ | – | – | 1 | – | 1 | 2 | 1 | 4 | 3 |
| N_{ev} | 18 | 7 | 21 | 2 | 10 | 1 | 1 | 9 | 3 |

Table 6. Charge-topology distribution of white stars in the dissociation of ^{16}O nuclei with an energy of 200 GeV per nucleon

| | | | | | | | | | | | |
|-----------|----|---|----|---|---|---|---|---|---|---|---|
| Z_f | 7 | 6 | 6 | 5 | 5 | 4 | 3 | 3 | – | – | – |
| $N_{Z=1}$ | 1 | – | 2 | 1 | 3 | 2 | 1 | 3 | – | 2 | 4 |
| $N_{Z=2}$ | – | 1 | – | 1 | – | 1 | 2 | 1 | 4 | 3 | 2 |
| N_{ev} | 49 | 6 | 10 | 5 | 1 | 3 | 2 | 2 | 2 | 4 | 2 |

Table 7. Charge-topology distribution of white stars in the dissociation of ^{10}B nuclei with an energy of 1 GeV per nucleon

| | | | | |
|-----------|---|---|---|----|
| Z_f | 4 | 3 | – | – |
| $N_{Z=1}$ | 1 | – | 3 | 1 |
| $N_{Z=2}$ | – | 1 | 1 | 2 |
| N_{ev} | 1 | 5 | 5 | 30 |

Table 8. Charge-topology distribution of white stars in the dissociation of ^{14}N nuclei with an energy of 2.1 GeV per nucleon

| | | | | | | | | |
|-----------|---|---|---|---|---|---|---|----|
| Z_f | 6 | 5 | 5 | 4 | 3 | 3 | – | – |
| $N_{Z=1}$ | 1 | – | 2 | 1 | 4 | 2 | 3 | 1 |
| $N_{Z=2}$ | – | 1 | – | 1 | – | 1 | 2 | 3 |
| N_{ev} | 6 | 3 | 2 | 1 | 1 | 1 | 1 | 10 |

Table 9. Charge-topology distribution of white stars in the dissociation of ^7Be nuclei of energy 1.23 GeV per nucleon

| | | | | |
|-----------|---|---|----|----|
| Z_f | 3 | – | – | – |
| $N_{Z=1}$ | 1 | 4 | 2 | – |
| $N_{Z=2}$ | – | – | 1 | 2 |
| N_{ev} | 7 | 2 | 38 | 28 |

Number of events of ^6Li coherent dissociation

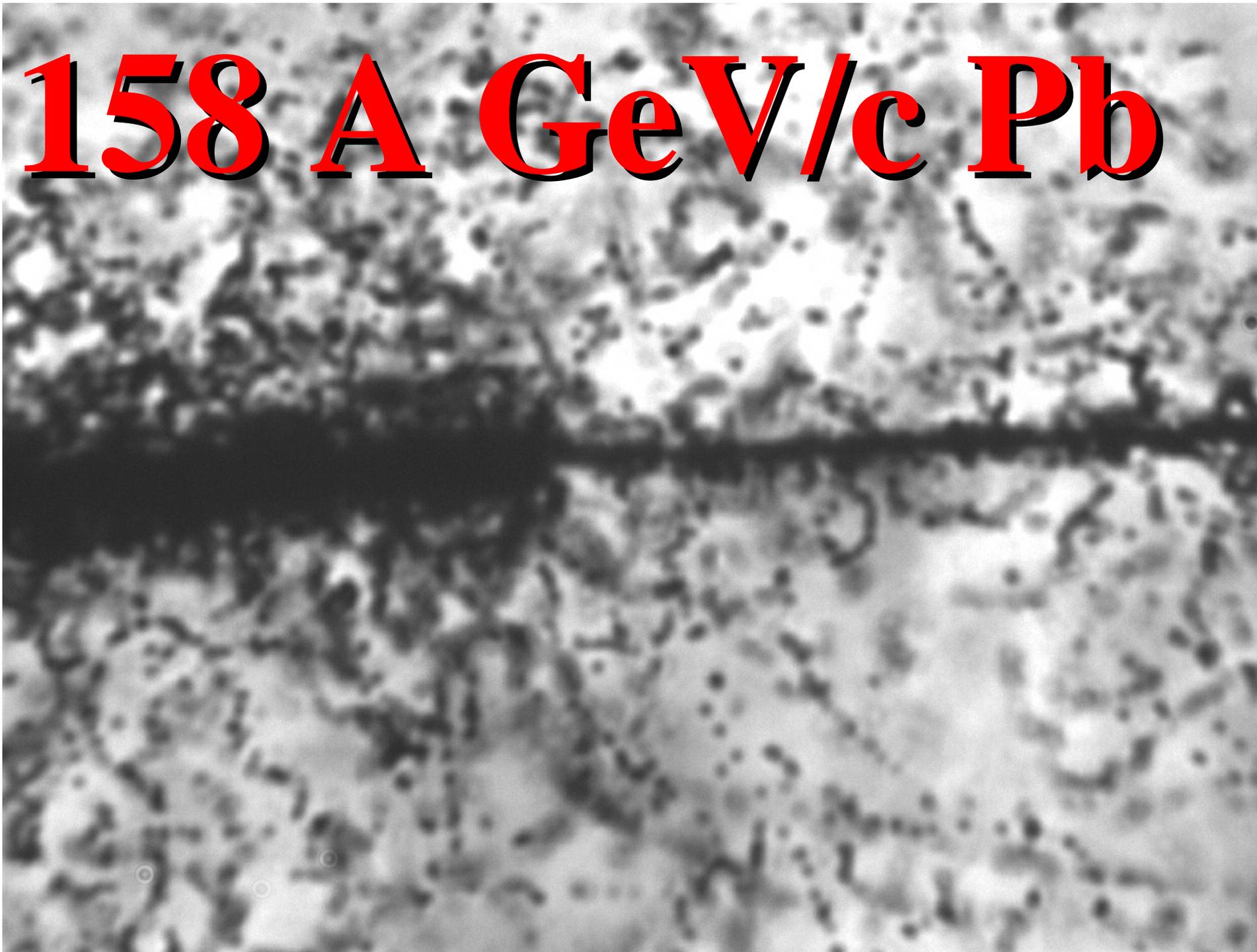
| Dissociation channel | Number of events | |
|----------------------|--|--|
| | without the excitation of the target nucleus ($N_h = 0$) | with the excitation of the target nucleus ($N_h \neq 0$) |
| $^4\text{He} + d$ | 23 | 24 |
| $^3\text{He} + t$ | 4 | 1 |
| $t + d + p$ | 4 | 3 |
| $d + d + d$ | 0 | 2 |

*The common topological feature for fragmentation of the **Ne, Mg, and Si nuclei** consists in a suppression of **binary splitting** to fragments with charges larger than 2.*

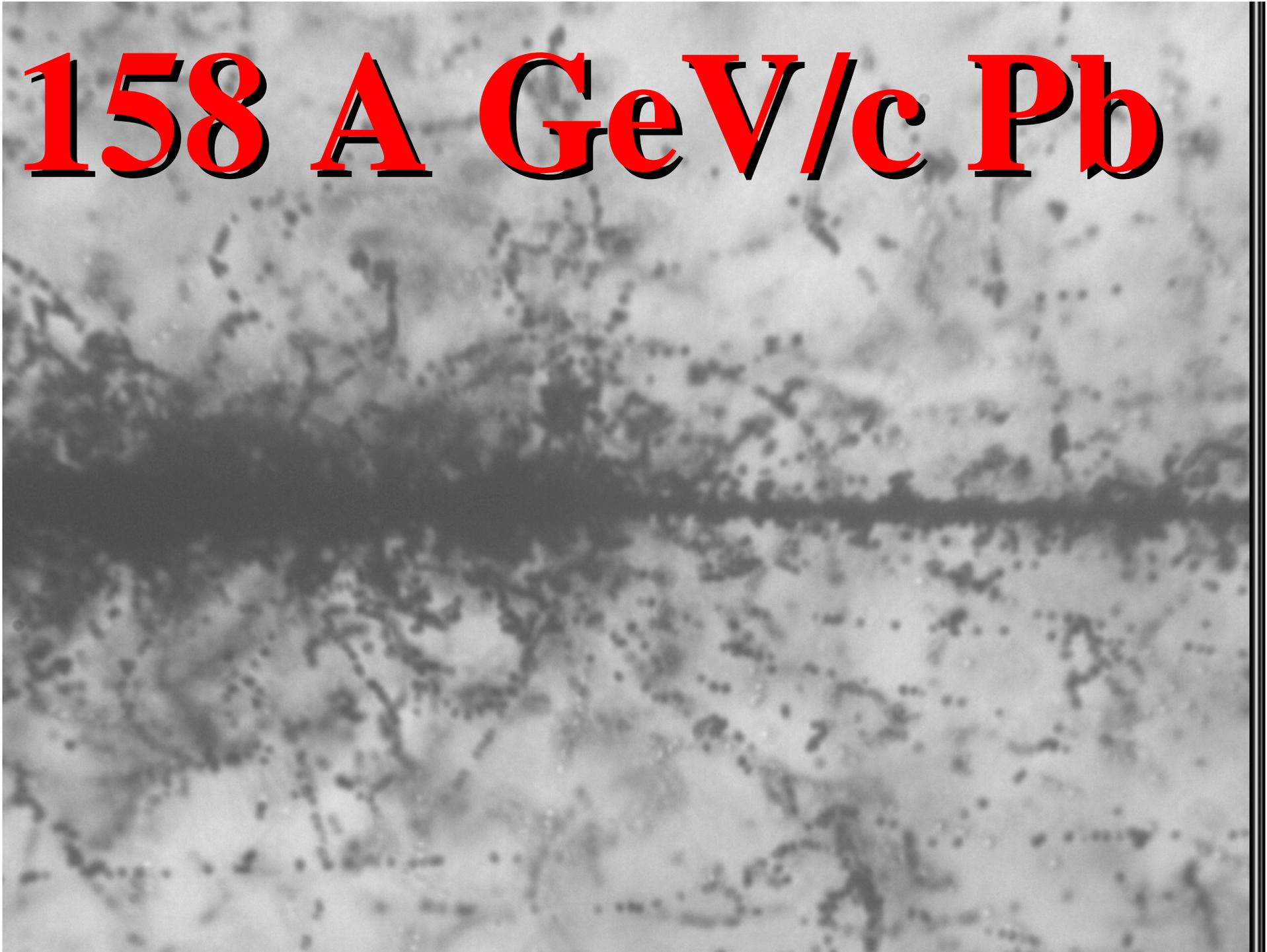
The growth of the fragmentation degree is revealed in an increase of the multiplicity of singly and doubly charged fragments up to complete dissociation with increasing of excitation.

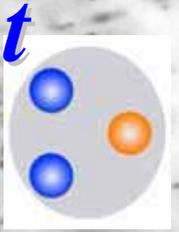
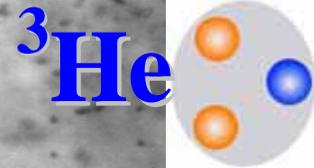
This circumstance shows in an obvious way on a domination of the multiple cluster states having high density over the binary states having lower energy thresholds.

158 A GeV/c Pb

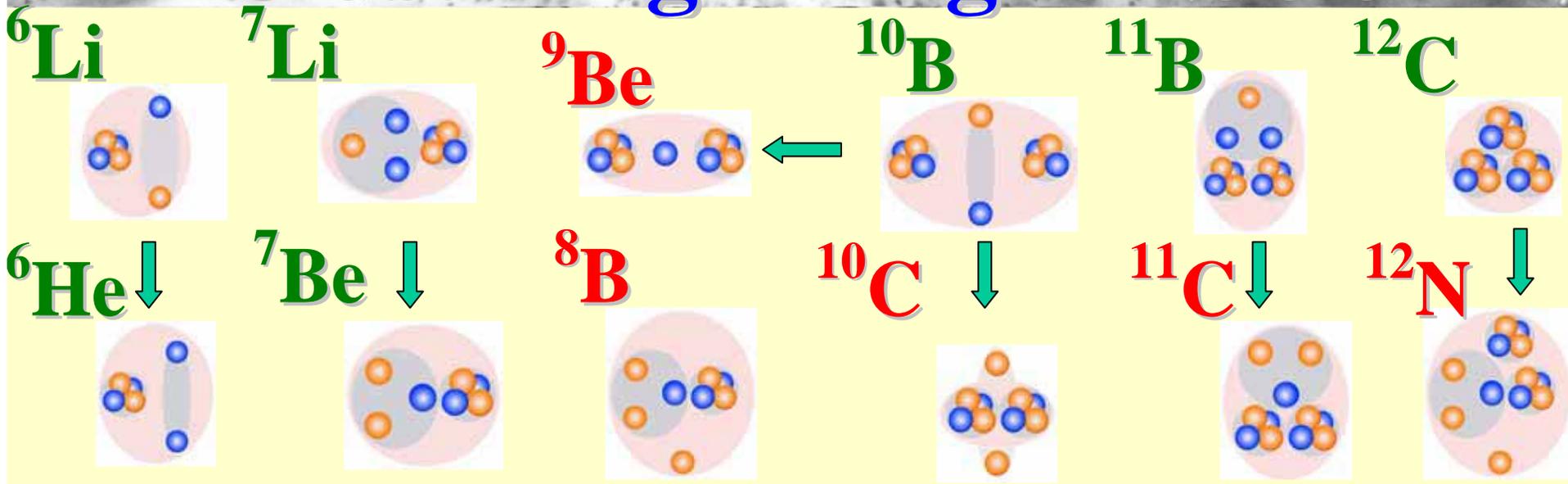


158 A GeV/c Pb





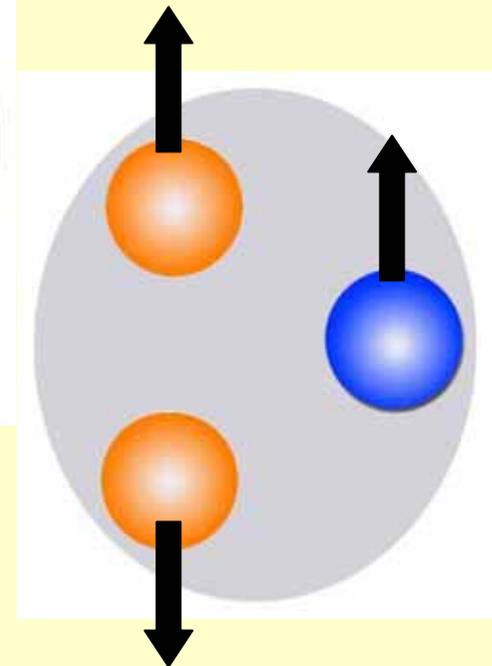
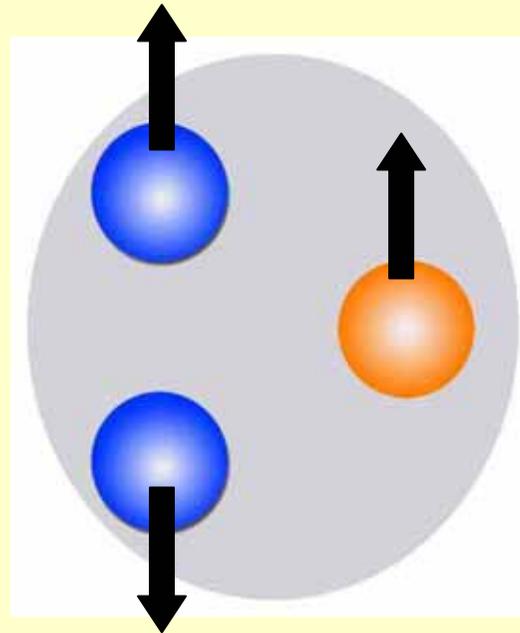
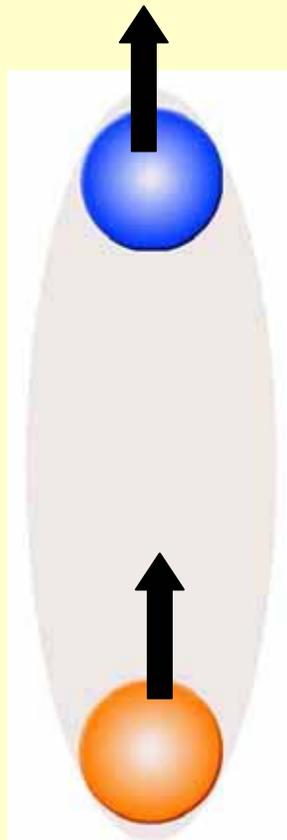
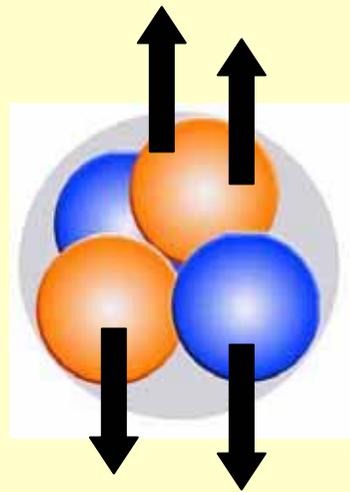
Clustering in Light Nuclei



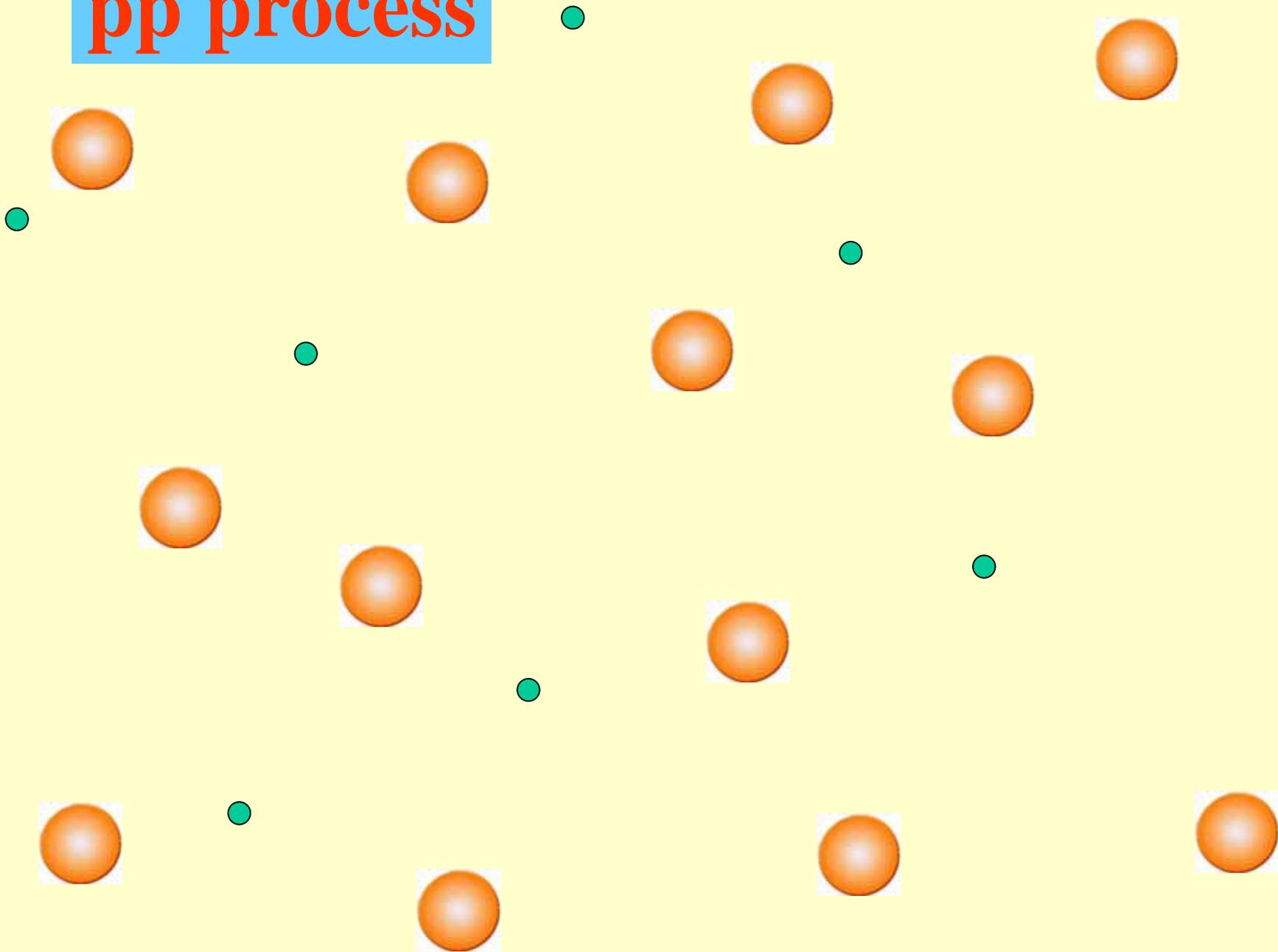
Secondary beams of light radioactive nuclei will be produced mostly via charge exchange reactions. ^8B and ^9Be beams has been formed via fragmentation of ^{10}B .

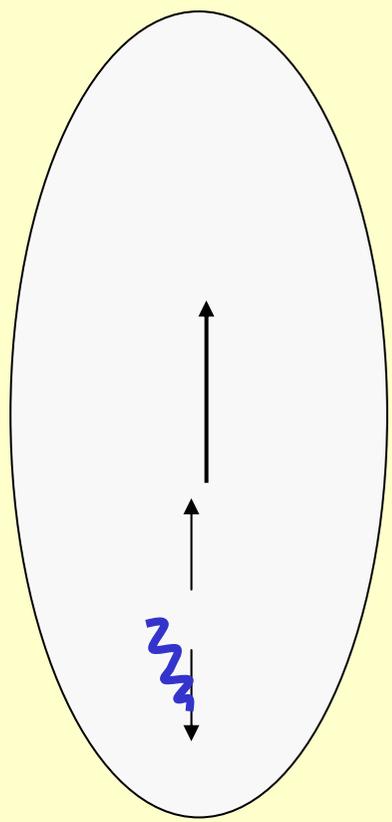
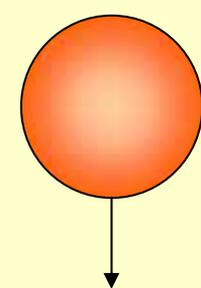
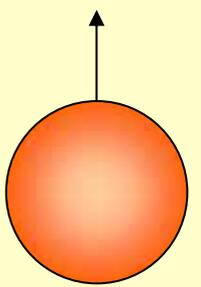
Clustering building blocks:

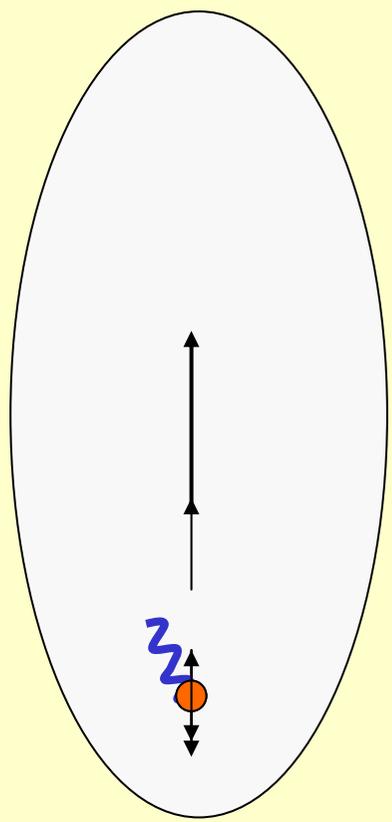
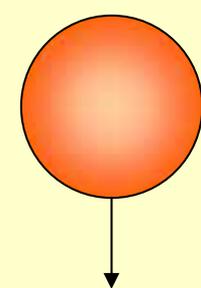
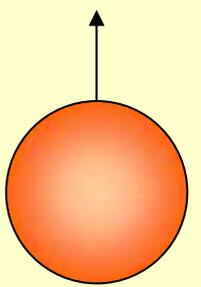
*more than one nucleon bound, stable & no excited states below particle decay thresholds –
deuteron, triton, ^4He , and ^3He nuclei*

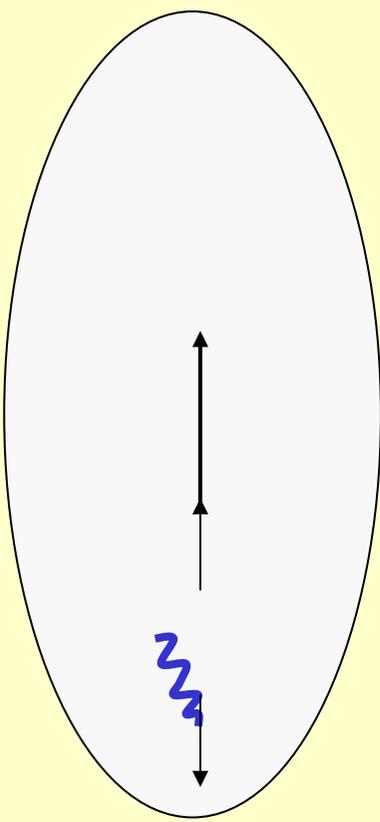
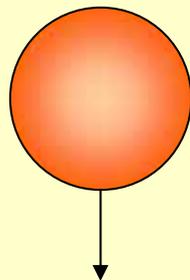
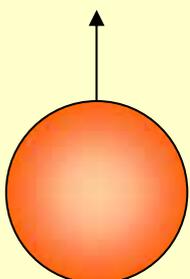


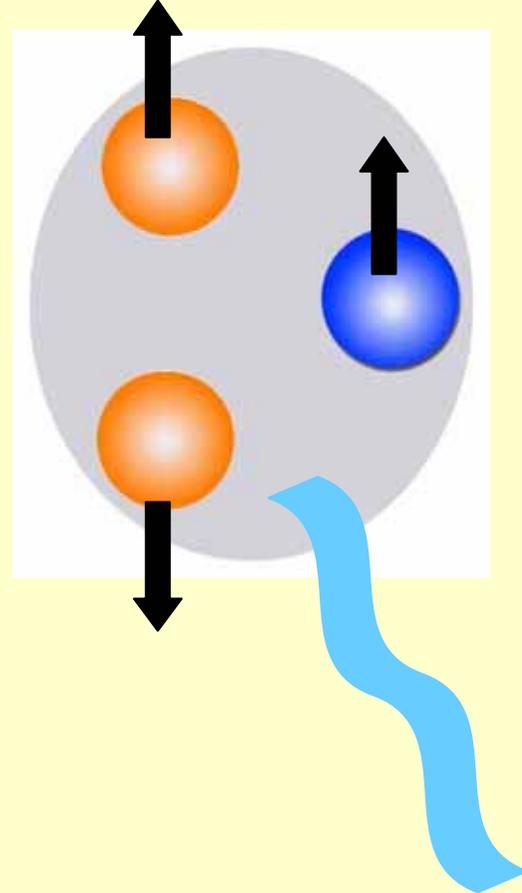
pp process



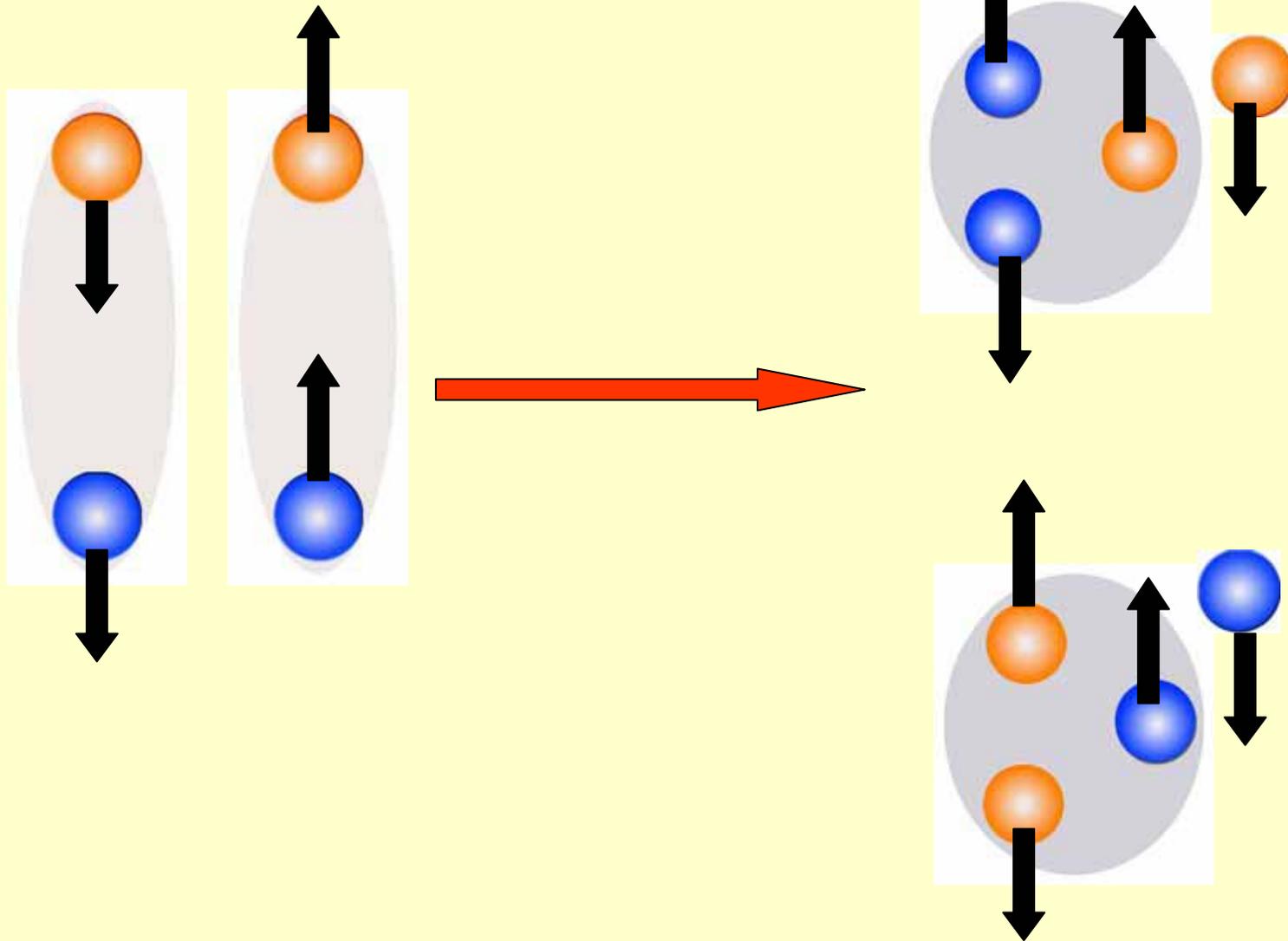




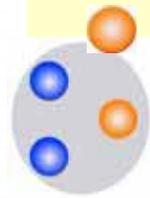
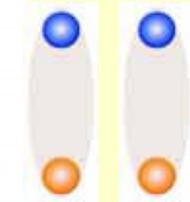
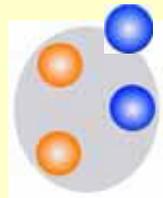
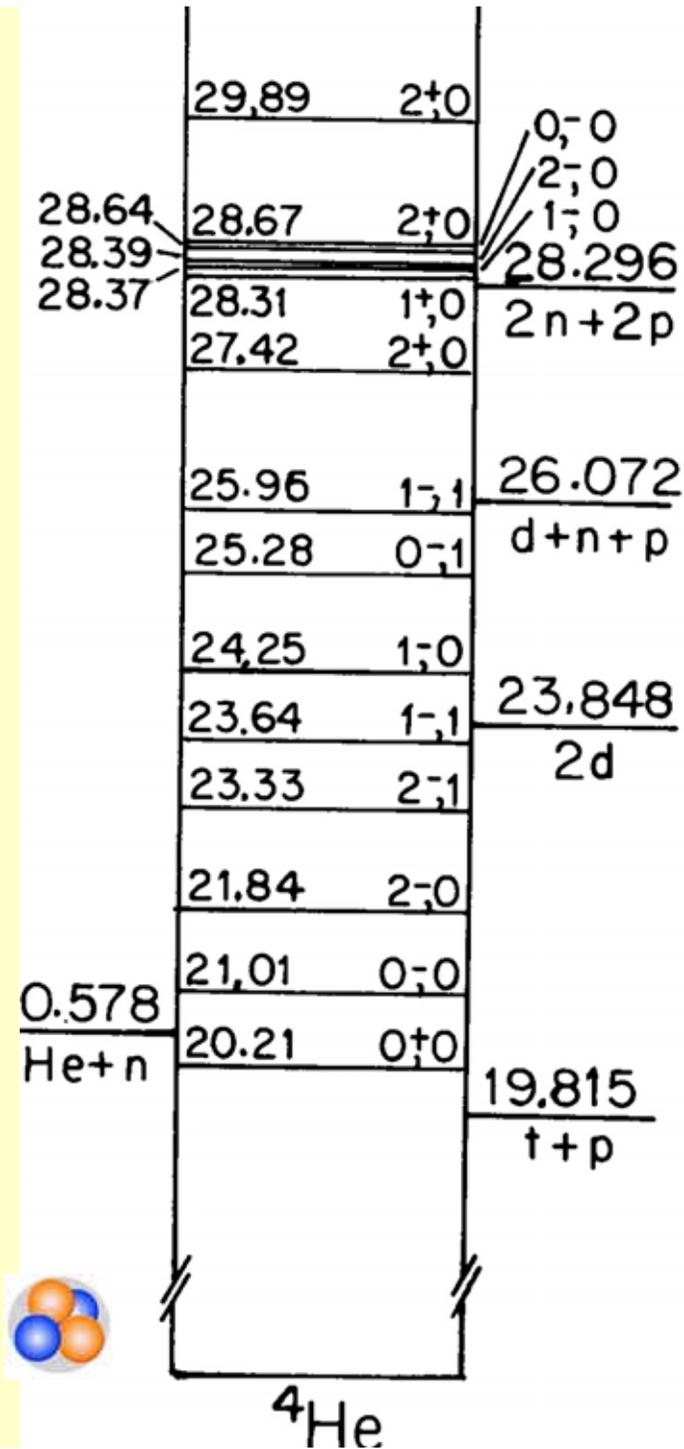




dd fusion

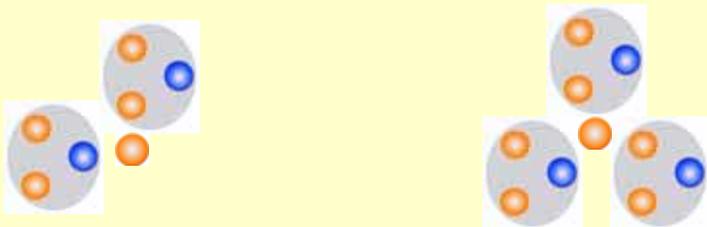
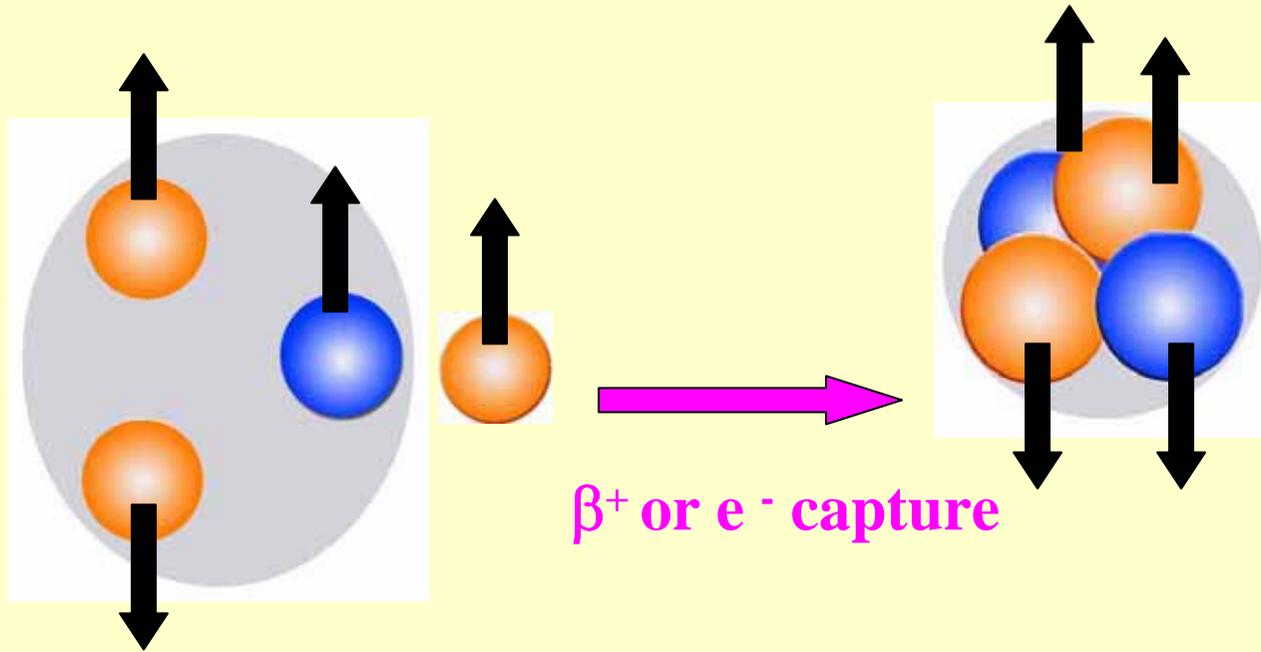


A = 4

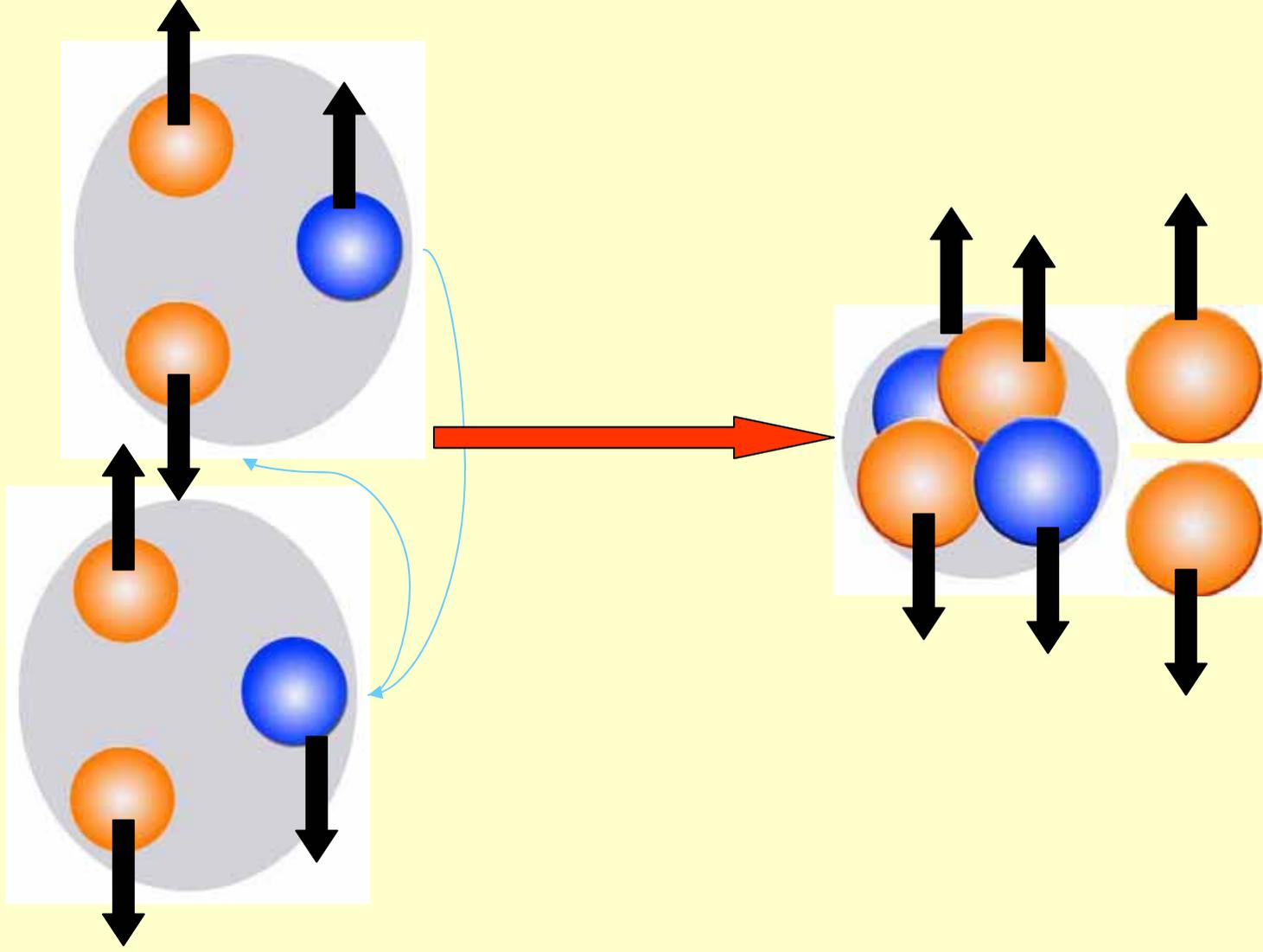


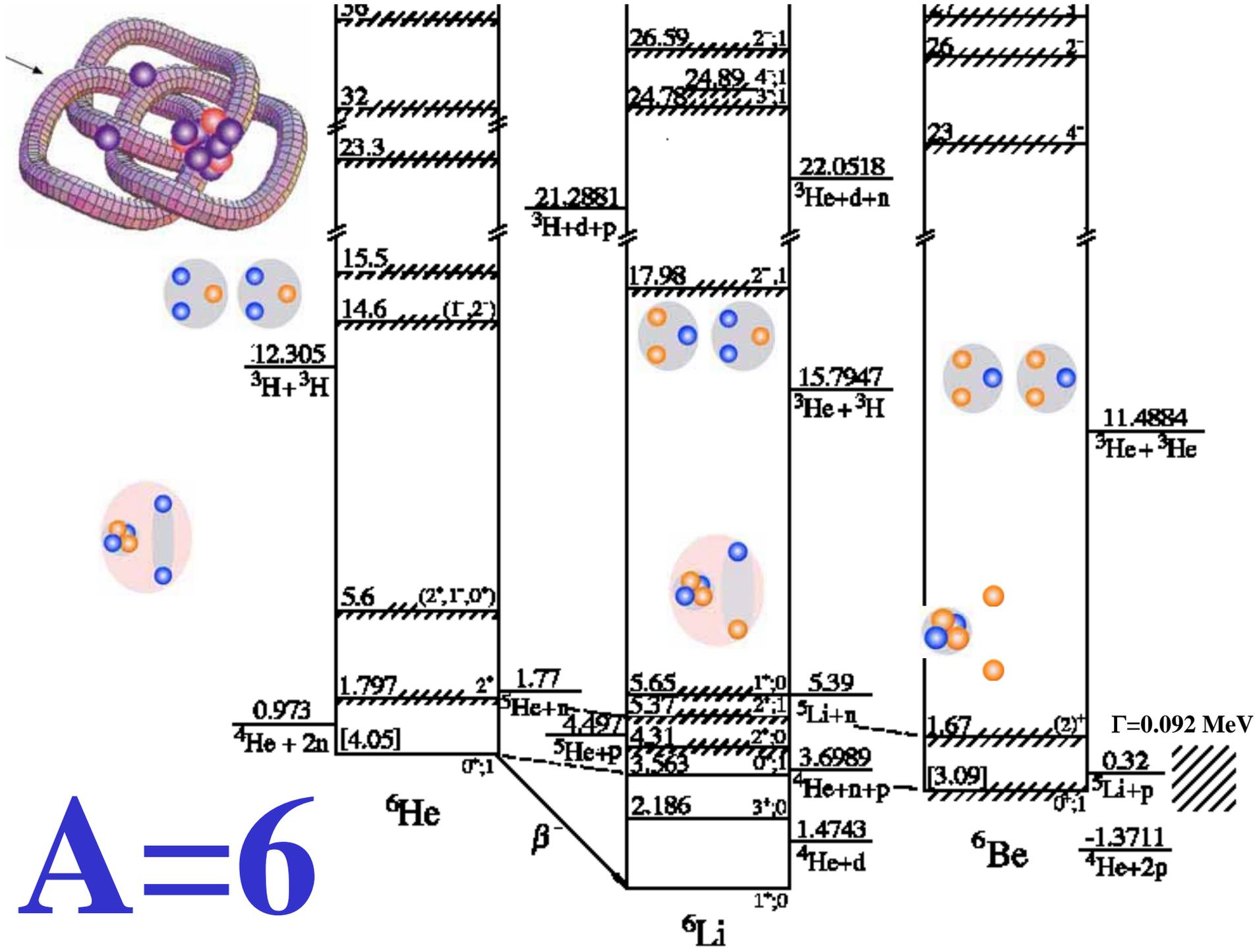
4He

Hep Process

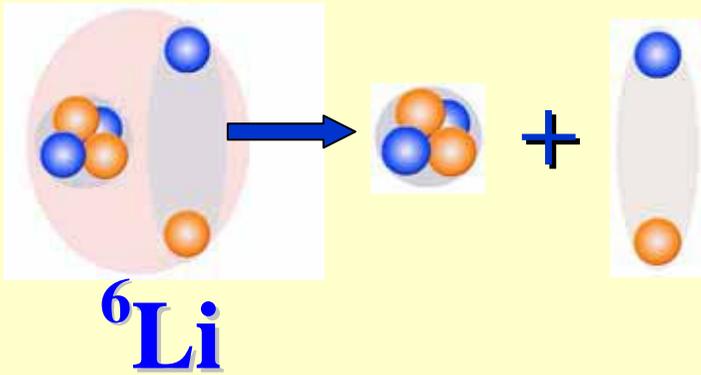


| | | |
|-----------------|-------|-------------------|
| | | 27.00 |
| 26.21 | 1^- | $n+3p$ |
| 25.44 | 0^- | |
| | | 24.78 |
| 23.68 | 1^- | $d+2p$ |
| 23.36 | 2^- | |
| ${}^4\text{Li}$ | | ${}^3\text{He}+p$ |
| | | 19.29 |

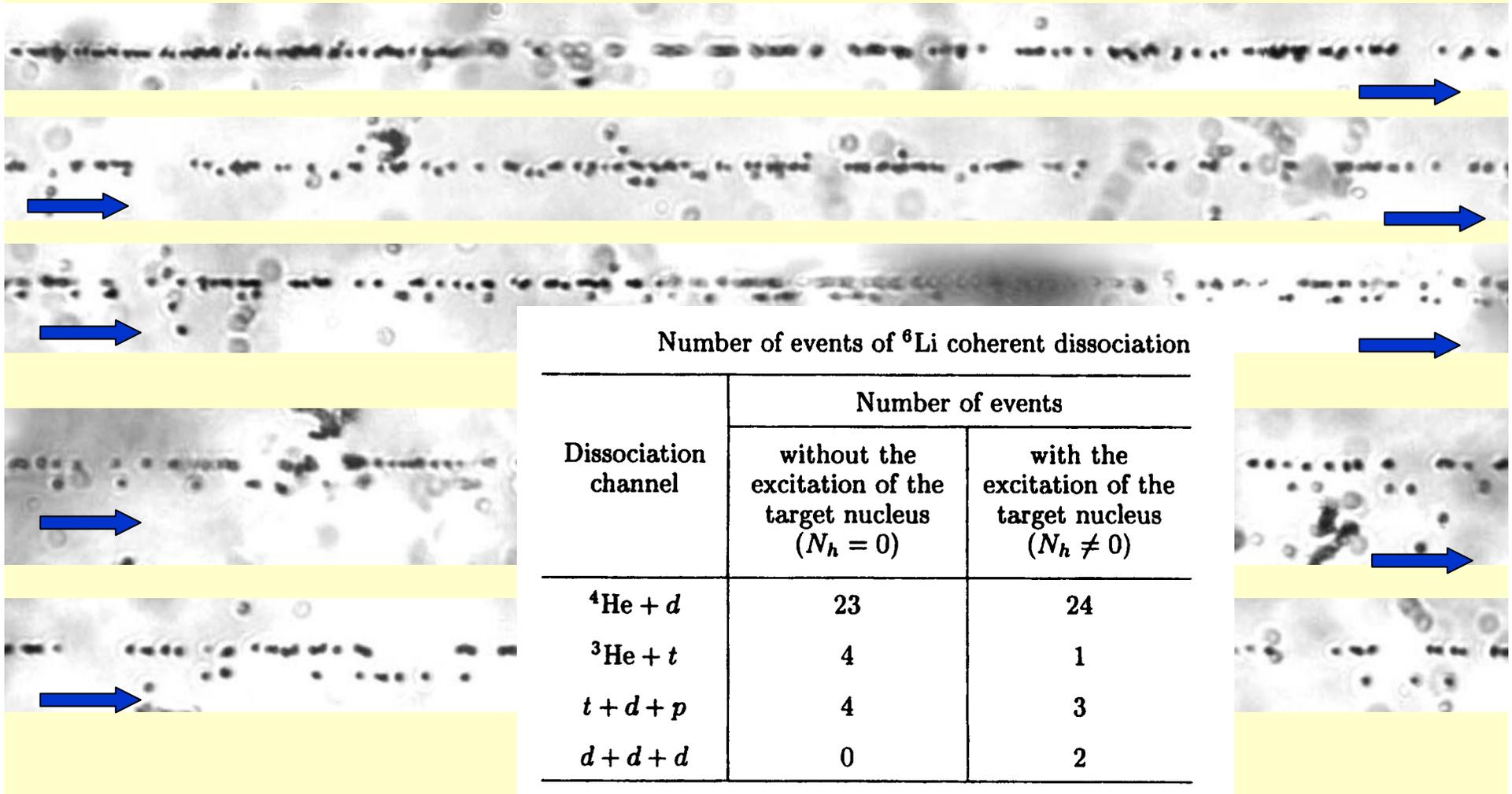




A=6



4.5A GeV/c ${}^6\text{Li}$ Coherent Dissociation (PAVICOM image)

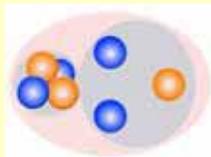


Number of events of ${}^6\text{Li}$ coherent dissociation

| Dissociation channel | Number of events | |
|----------------------|--|--|
| | without the excitation of the target nucleus ($N_h = 0$) | with the excitation of the target nucleus ($N_h \neq 0$) |
| ${}^4\text{He} + d$ | 23 | 24 |
| ${}^3\text{He} + t$ | 4 | 1 |
| $t + d + p$ | 4 | 3 |
| $d + d + d$ | 0 | 2 |

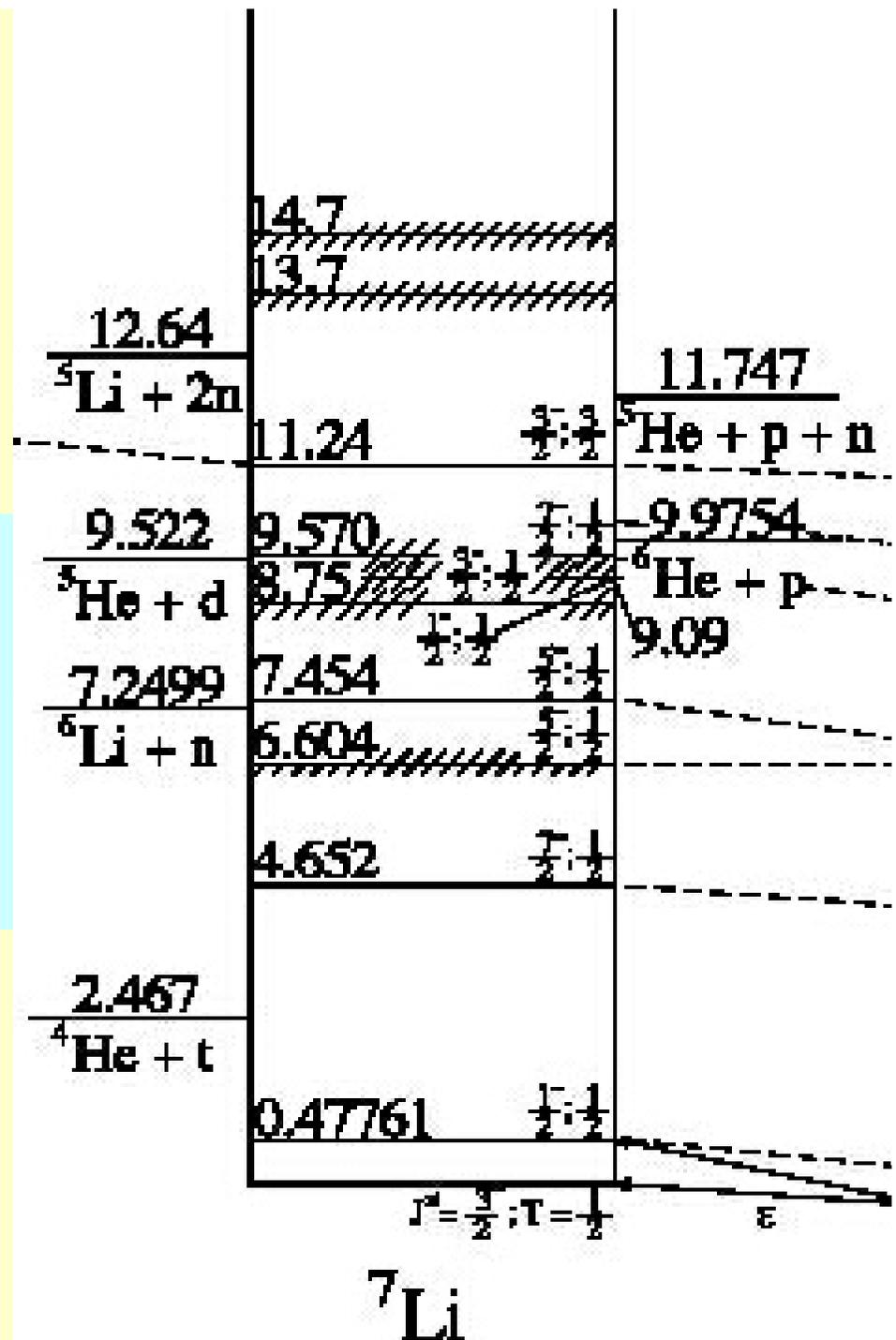
Triton Clustering

${}^7\text{Li}$

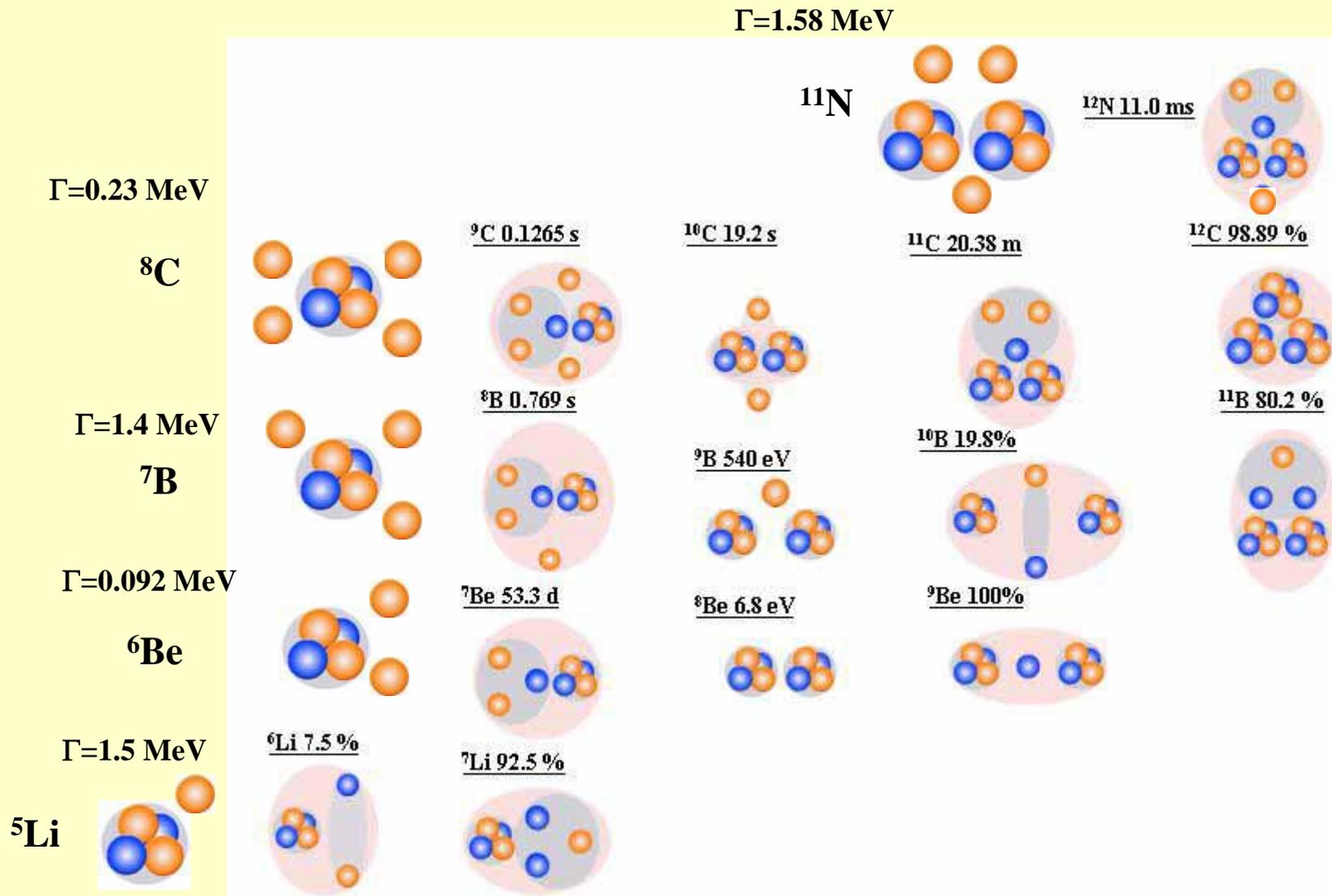


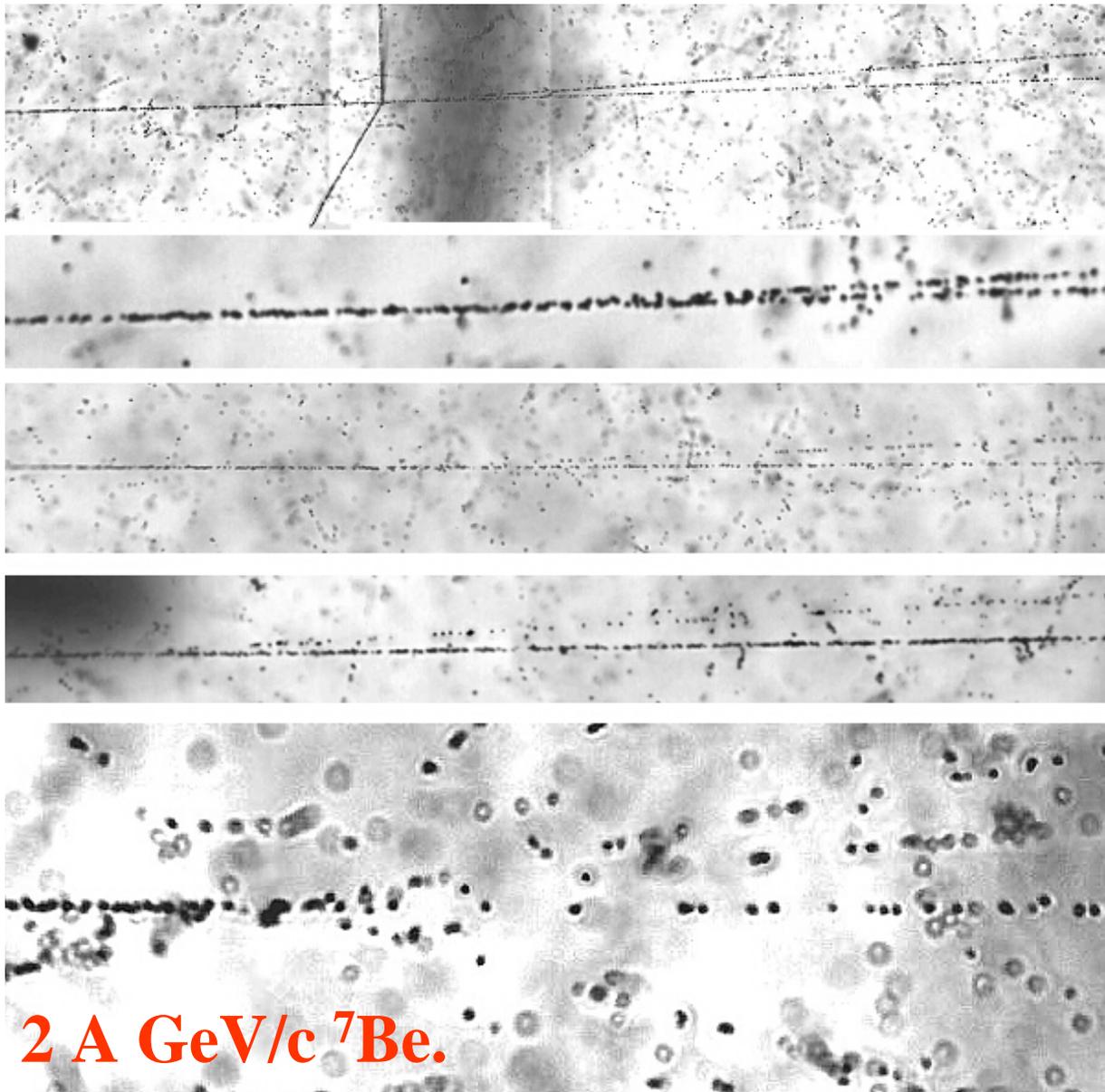
${}^7\text{Li}$. About 7% of all inelastic interactions of ${}^7\text{Li}$ nuclei are “white” stars (80 events).

Decay of ${}^7\text{Li}$ nucleus to α -particle and triton - 40 events.

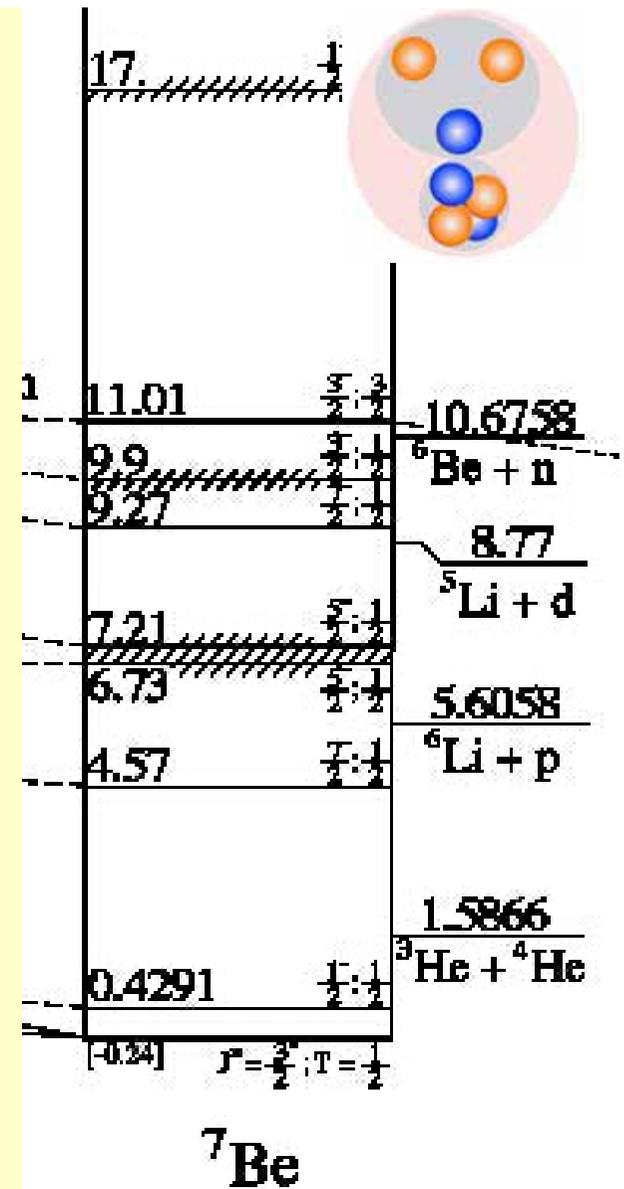


Toward stability frontier



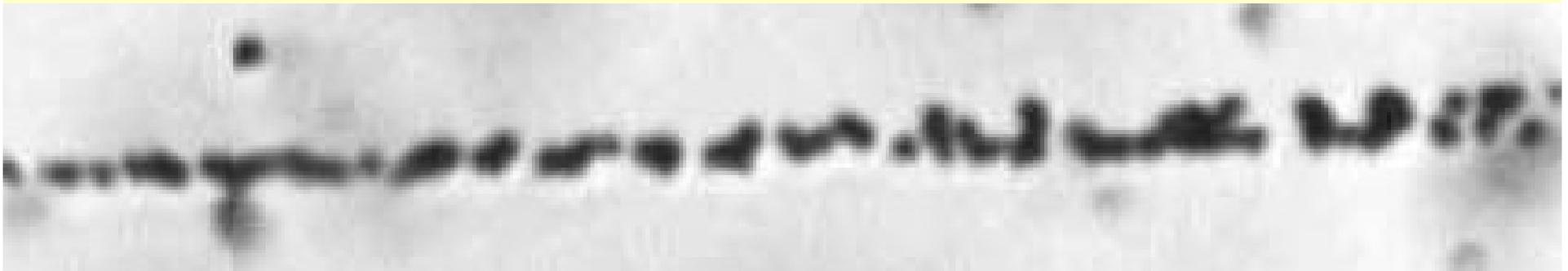
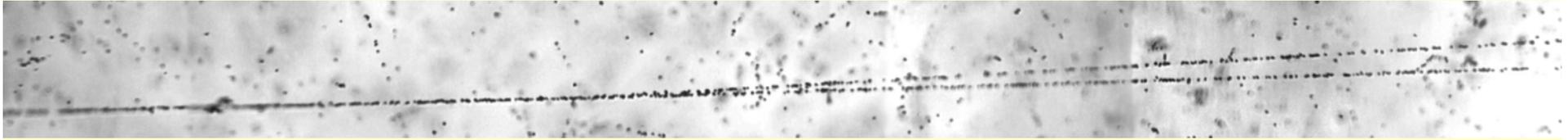


2 A GeV/c ${}^7\text{Be}$.



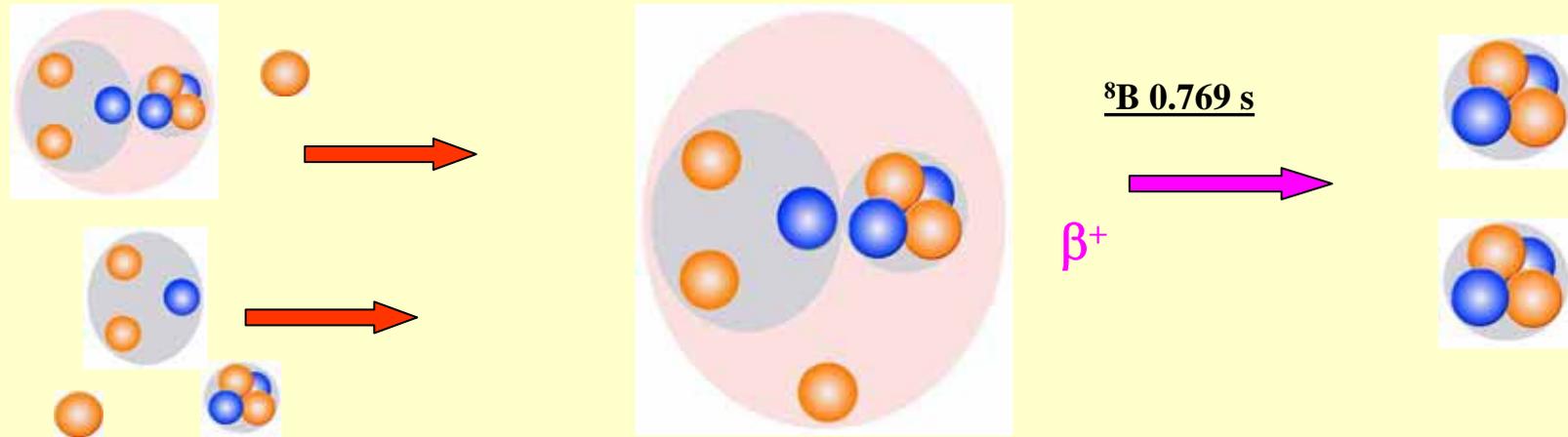
*Splitting to HeHe with two target fragments,
HeHe, HeHH, ${}^6\text{Li}p$, and 4H.*

Relativistic ${}^7\text{Be}$ fragmentation: 2+2



The ${}^7\text{Be}^* \rightarrow \alpha {}^3\text{He}$ decay is observed in 22 “white stars” with 2+2 topology. In the latter, 5 “white” stars are identified as the ${}^7\text{Be}^* \rightarrow (n) {}^3\text{He} {}^3\text{He}$ decay. Thus, a ${}^3\text{He}$ clustering is clearly demonstrated in dissociation of the ${}^7\text{Be}$ nucleus.

“Ternary H&He Process”

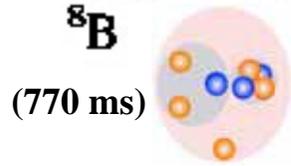
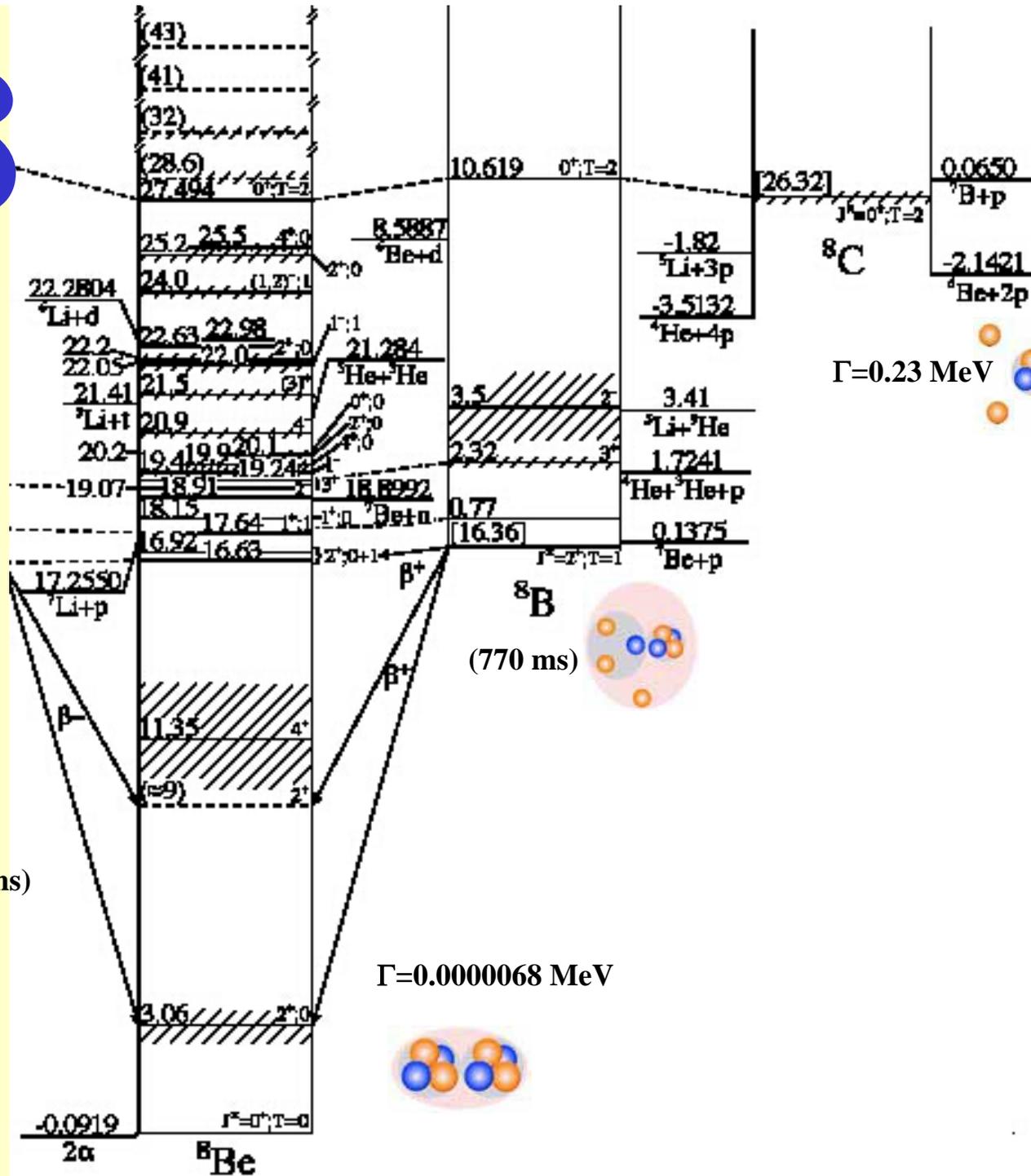
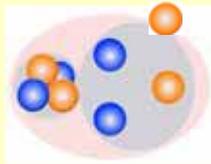
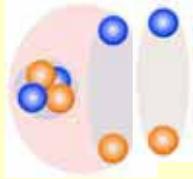


The ^{10}B nuclei with a momentum of $2A\text{ GeV}/c$ and an intensity of about 10^8 nuclei per cycle were accelerated at the JINR nuclotron. A beam of secondary nuclei of a magnetic rigidity corresponding to $Z/A = 5/8$ ($^{10}\text{B} \rightarrow ^8\text{B}$ fragmentation) was provided for emulsions.

We plan to determine the probabilities

$^8\text{B} \rightarrow ^7\text{Bep}$ (9), $^3,^4\text{He}^3\text{Hep}$ (8), $^6\text{Lipp}$ (1),
 HeHHp (7), and HHHpp (1).

A = 8

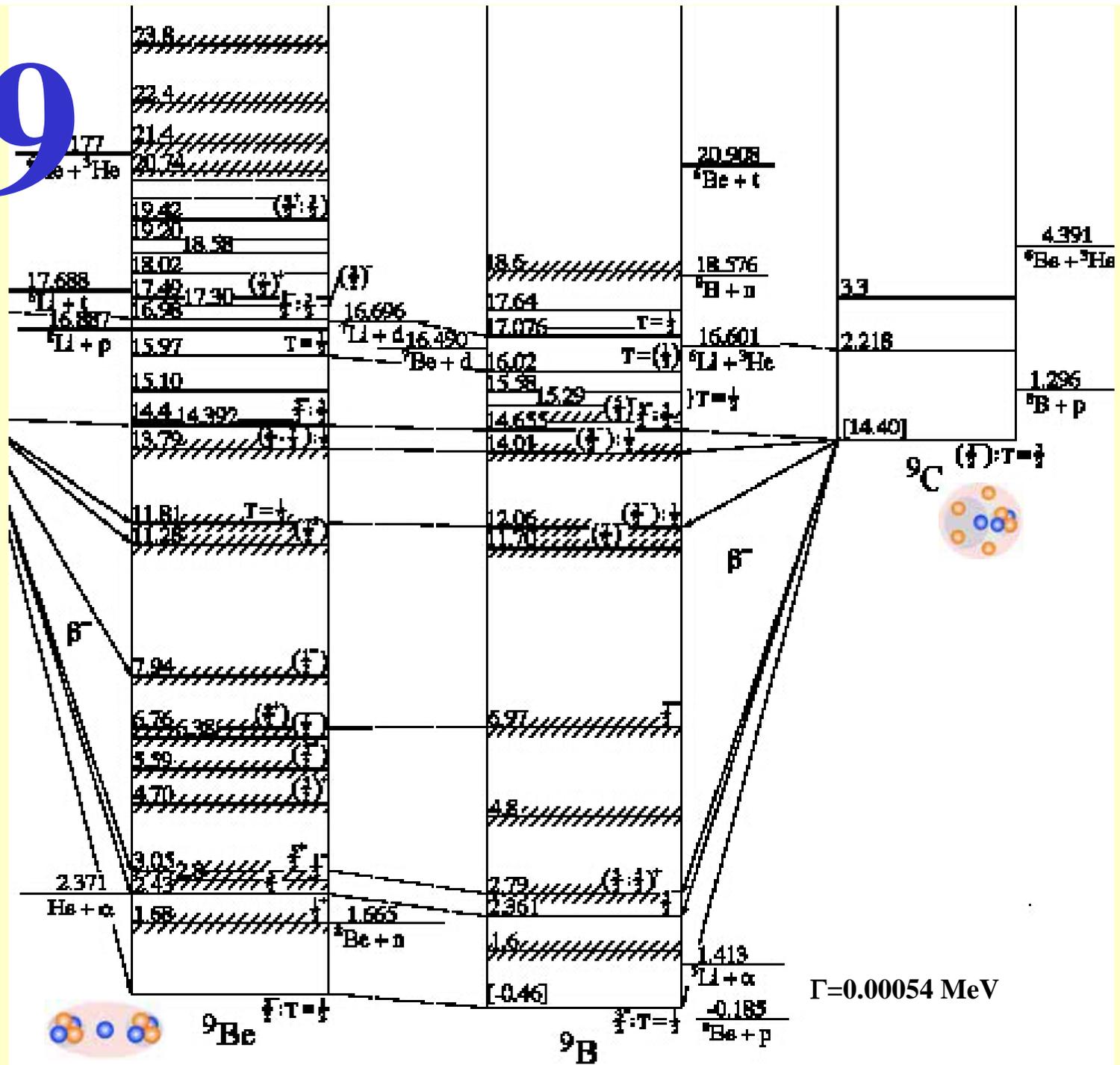


$\Gamma=0.0000068 \text{ MeV}$



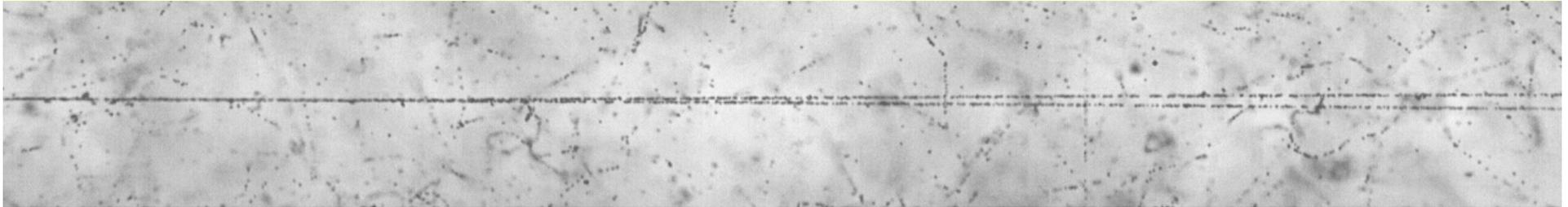
${}^8\text{He}$ (118 ms), ${}^8\text{Li}$ (838 ms)

A=9

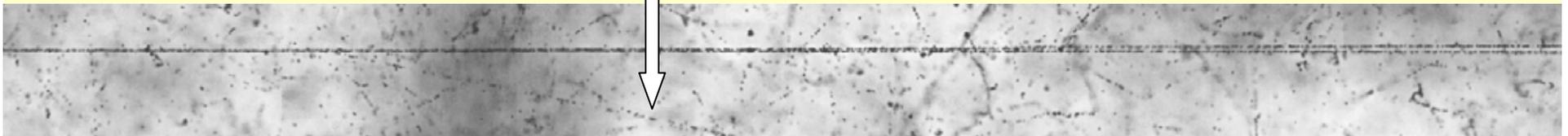


1.3A GeV ^9Be dissociation in 2+2. $^{10}\text{B} \rightarrow ^9\text{Be}$, Nuclotron, 2004.

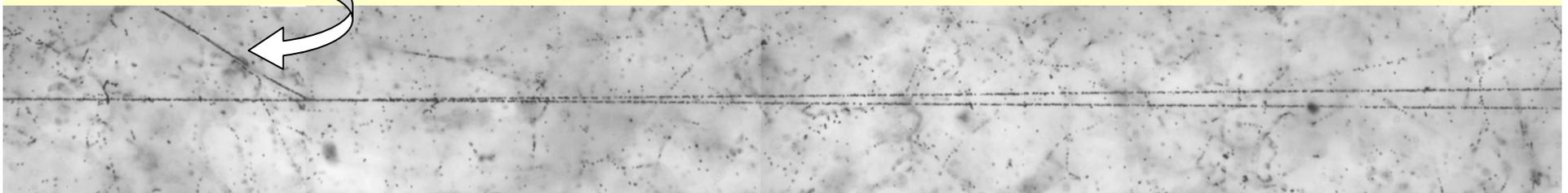
“white” star



with recoil proton



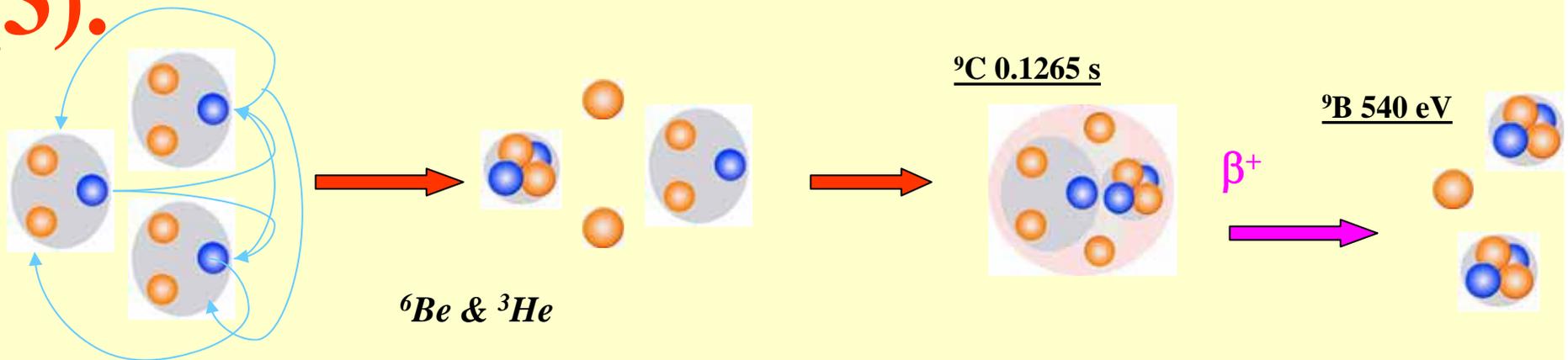
with heavy fragment of target nucleus

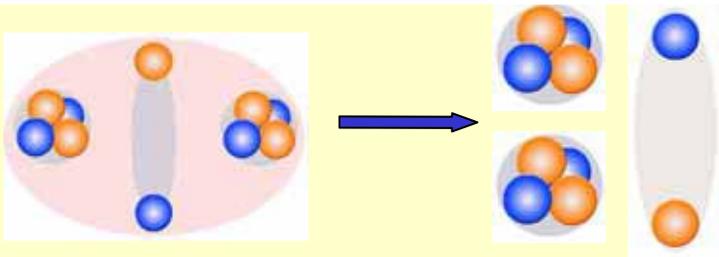


“Ternary ^3He Process”

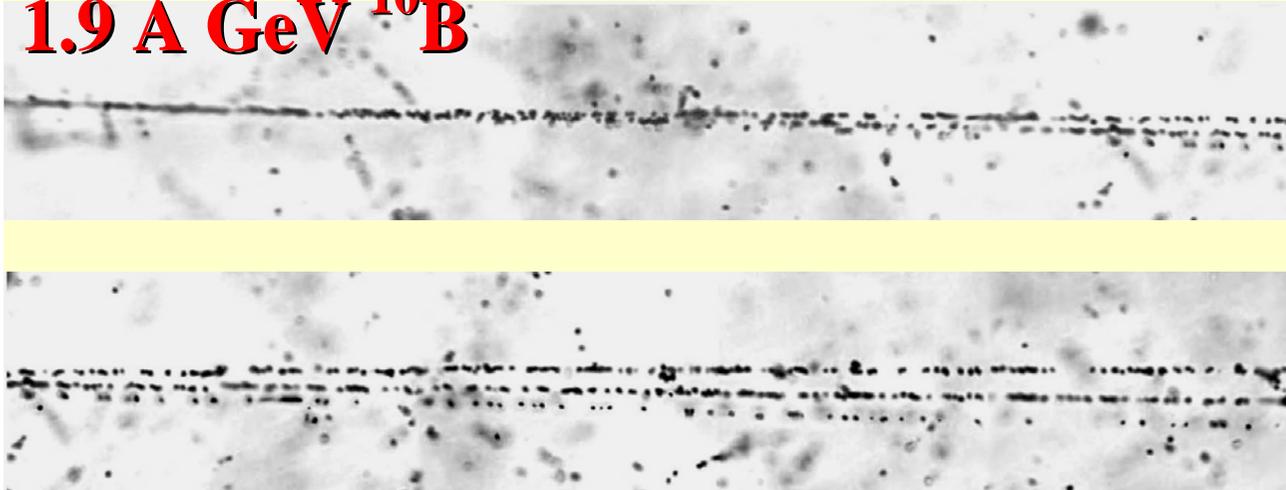
2 A GeV/c carbon beam of a magnetic rigidity $Z/A = 6/9$ ($^{12}\text{C} \rightarrow ^9\text{C}$) was provided for emulsions to determine the probabilities

$^9\text{C} \rightarrow ^8\text{Bp}$ (1) , $^7\text{Bepp}$ (2) , HeHepp (7), ^6Li pp (0) , HeHHpp (5) , HeHeHe (3).





1.9 A GeV ^{10}B

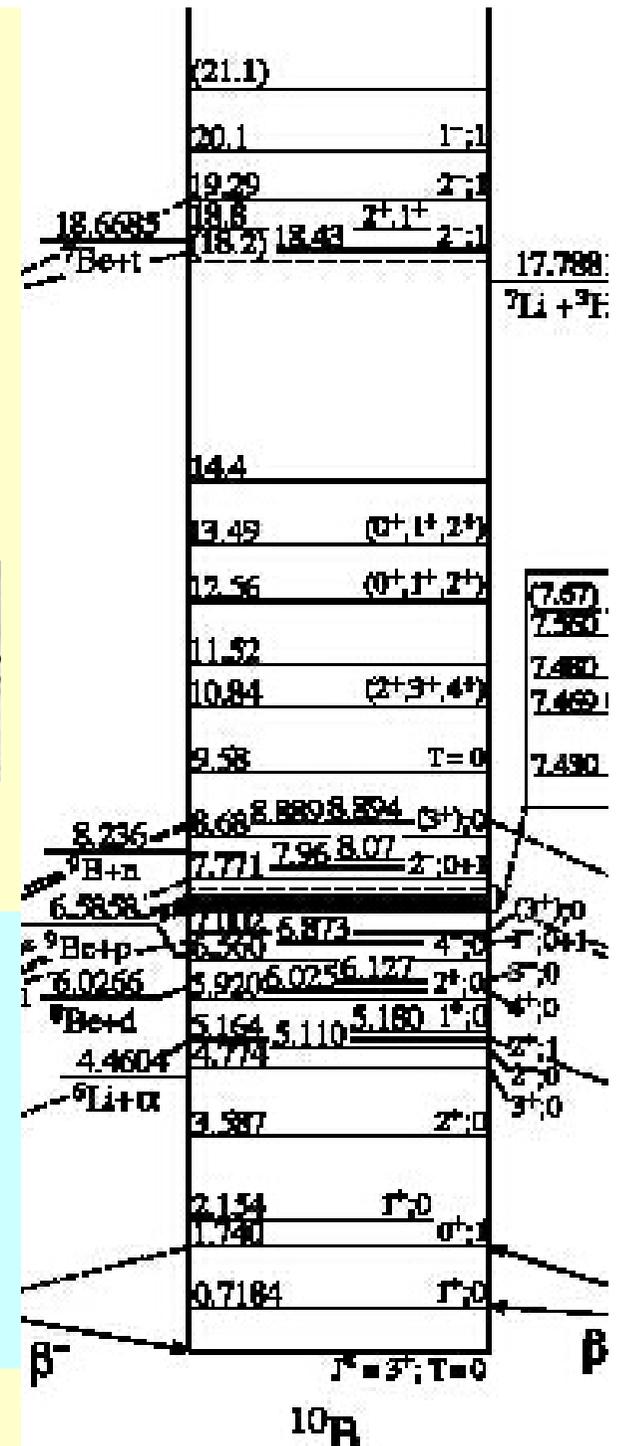


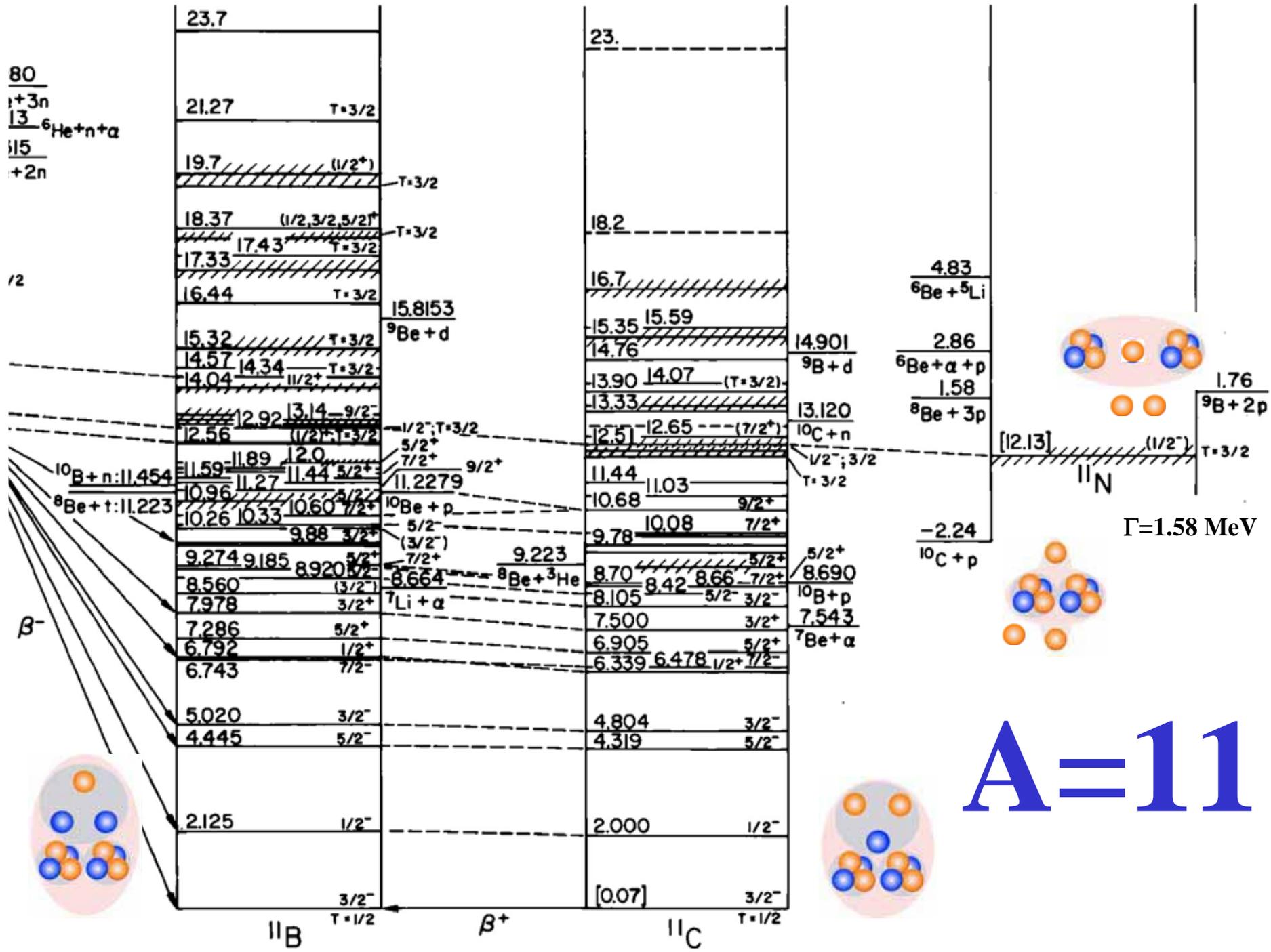
^{10}B is disintegrated to 2 doubly charged and 1 singly charged particles in 70% of “white” stars.

A singly charged particle is the deuteron in 40% like in case of ^6Li .

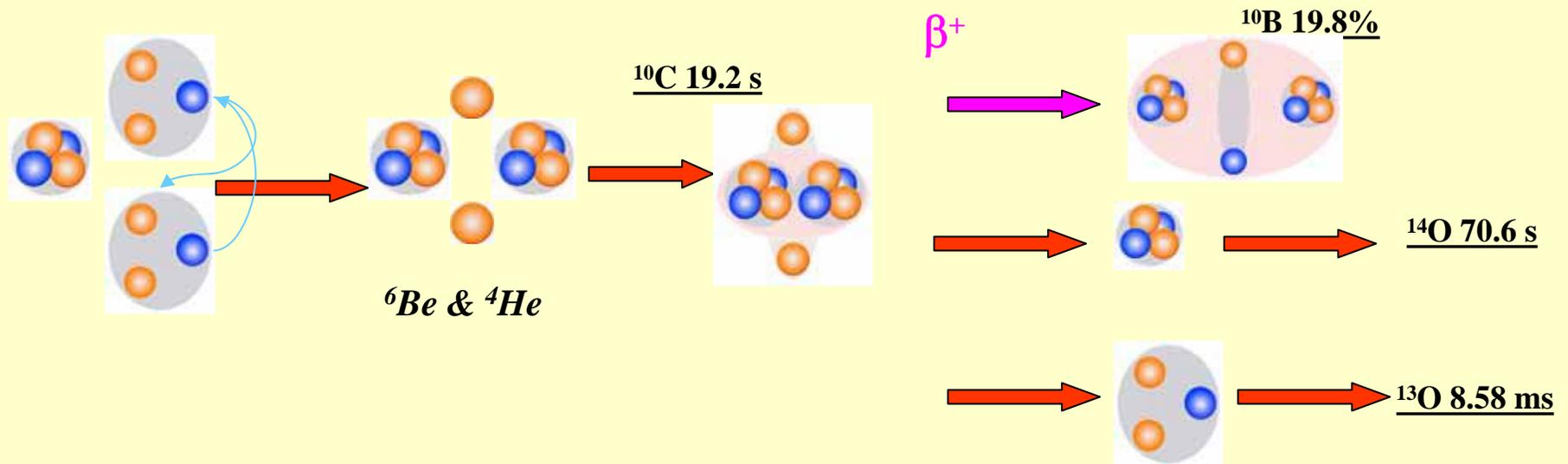
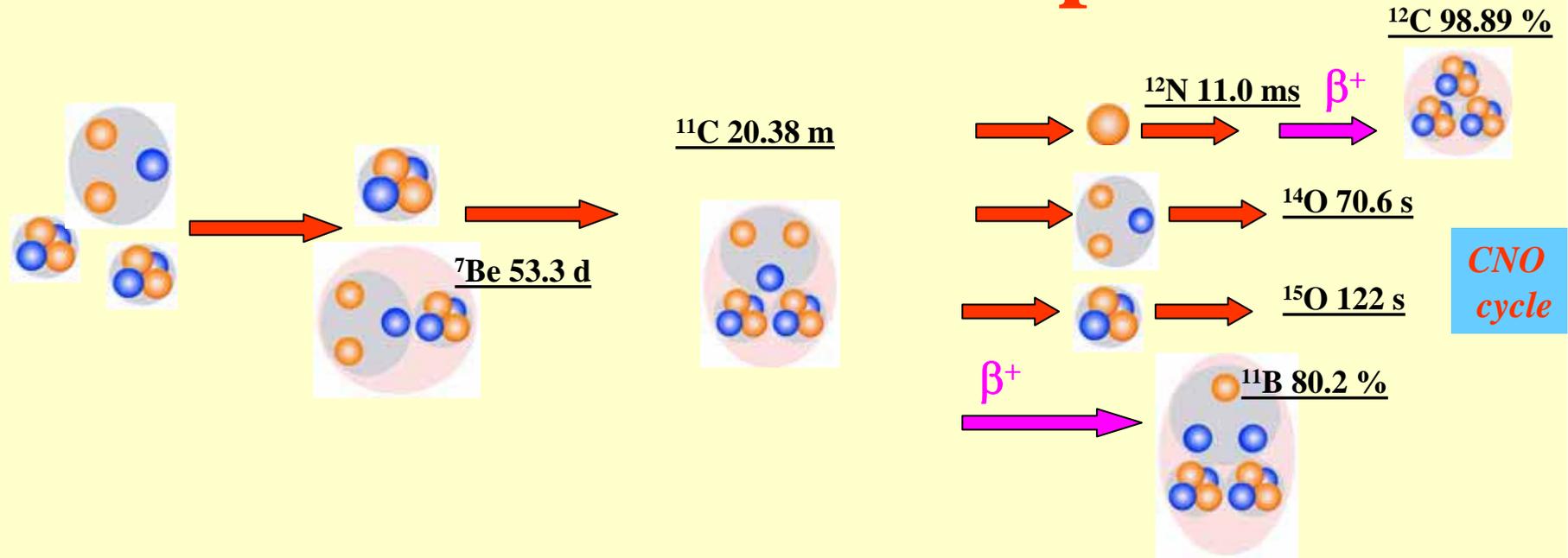
^8Be contribution is 20%.

$^{10}\text{B} \rightarrow ^9\text{Be} + p - 3\%$

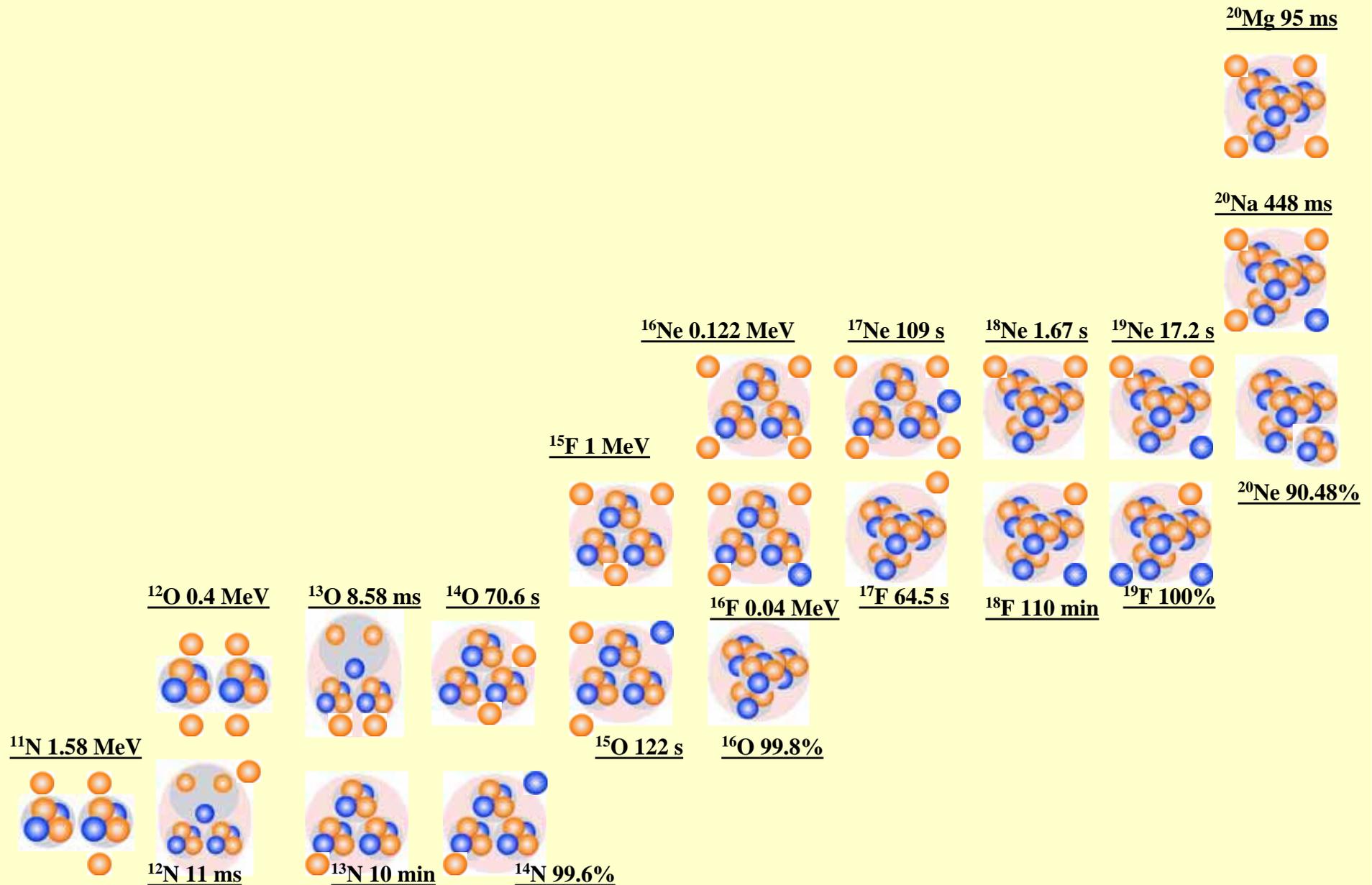




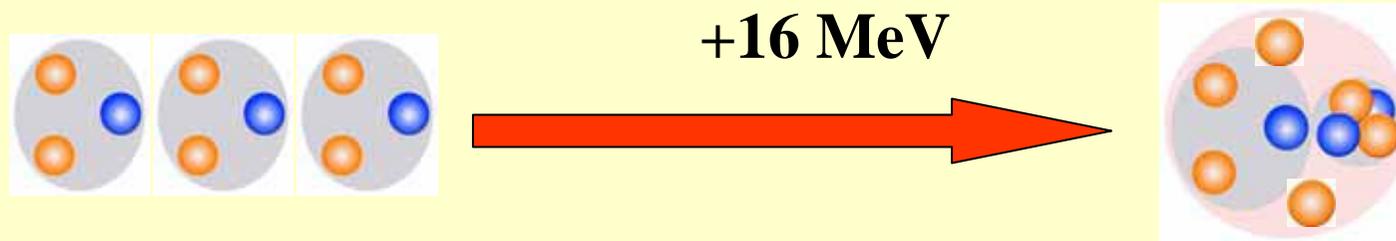
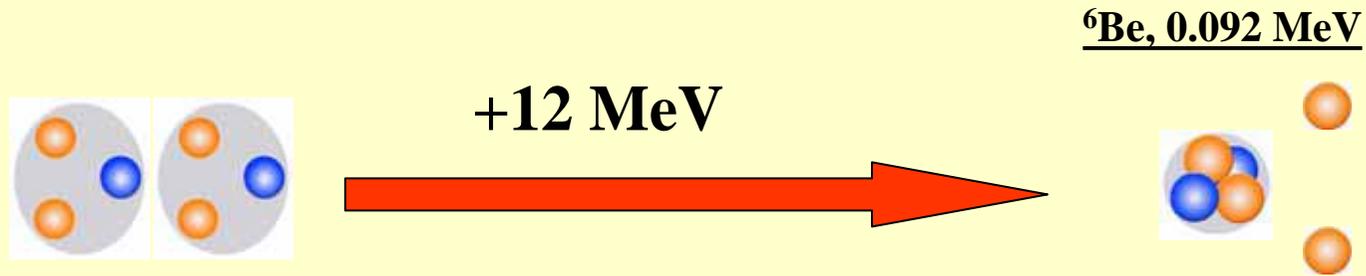
“ ^3He Process: mixed isotope fusion”



Walking along proton stability line

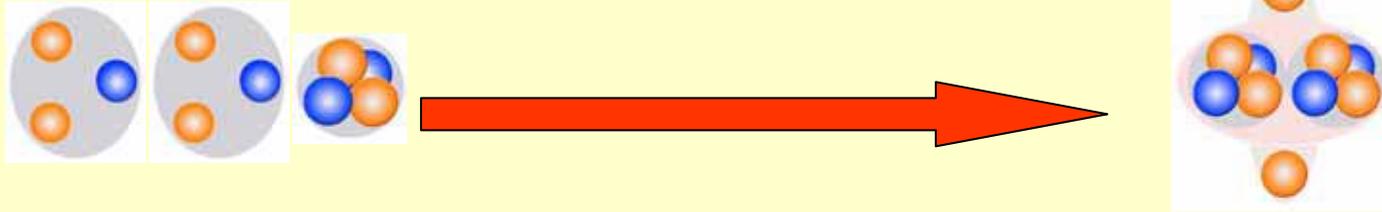


Multifragmentation in H&He



Multifragmentation in H&He

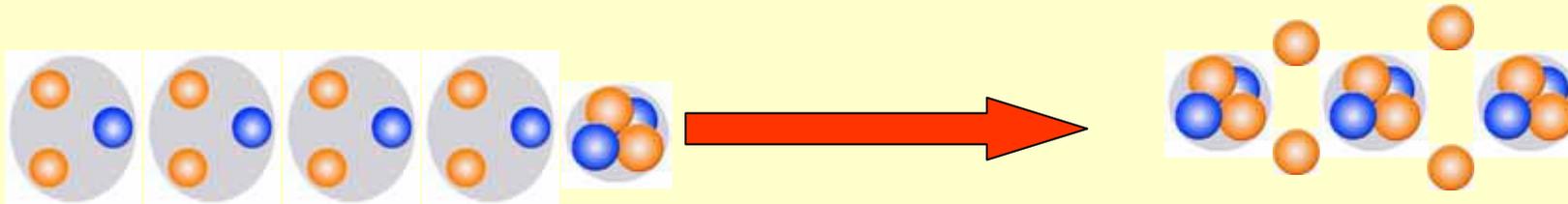
^{10}C 19.2 s



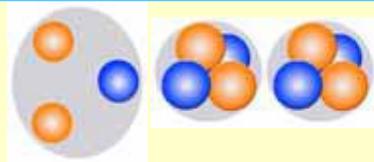
^{13}O 8.58 ms



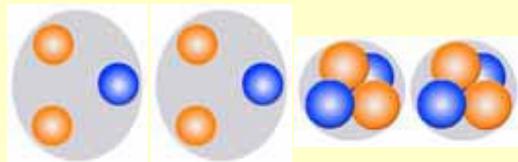
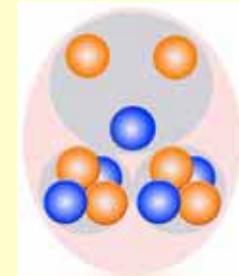
^{16}Ne 0.122 MeV



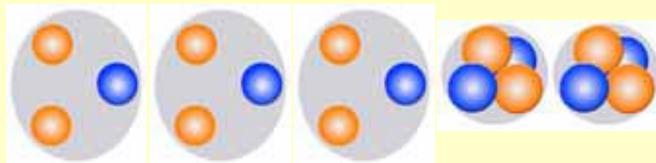
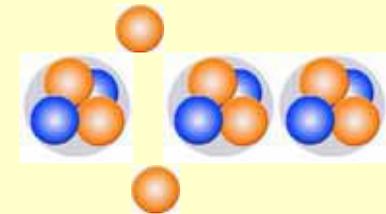
Multifragmentation in H&He



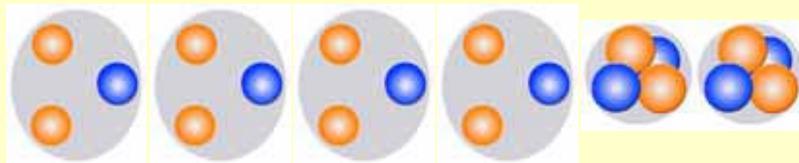
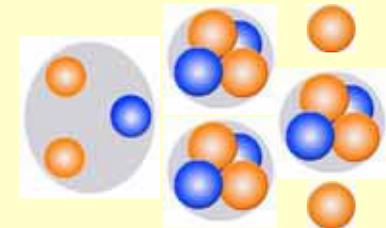
^{11}C 20.38 m



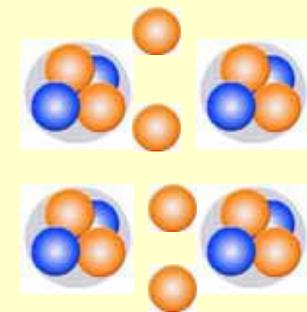
^{14}O 70.6 s



^{17}Ne 109 s



^{20}Mg 95 ms

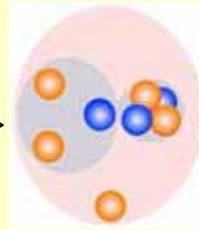
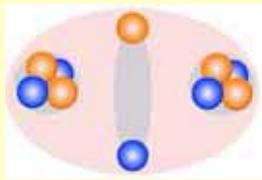


Fragmentation of relativistic nuclei provides an excellent quantum “laboratory” to explore the transition of nuclei from the ground state to a gas-like phase composed of nucleons and few-nucleon clusters having no excited states, i. e. d , t , ${}^3\text{He}$, and α .

The research challenge is to find indications for the formation of quasi-stable systems significantly exceeding the sizes of the fragments.

Search for such states is of interest since they can play a role of intermediate states for a stellar nuclear fusion due to dramatically reduced Coulomb repulsion.

The fragmentation features might assist one to disclose the scenarios of few-body fusions as processes inverse to fragmentation.



$2 A \text{ GeV}/c \text{ } ^{10}\text{B} \rightarrow \text{}^8\text{B}$

