Chapter 24 The Hoyle State in Relativistic ¹²C Dissociation



D. A. Artemenkov, M. Haiduc, N. K. Kornegrutsa, E. Mitsova, N. G. Peresadko, V. V. Rusakova, R. Stanoeva, A. A. Zaitsev, P. I. Zarubin and I. G. Zarubina

Abstract Production of α -particle triples in the Hoyle state (HS) in dissociation of ¹²C nuclei at 3.65 and 0.42 *A* GeV in nuclear track emulsion is revealed by the invariant mass approach. Contribution of the HS to the dissociation ¹²C \rightarrow 3 α is $(11 \pm 3)\%$. Reanalysis of data on coherent dissociation ¹⁶O \rightarrow 4 α at 3.65 *A* GeV is revealed the HS contribution of $(22 \pm 2)\%$.

Events of dissociation of relativistic light nuclei observable in detail in the nuclear track emulsion (NTE) contain holistic information on ensembles of lightest nuclei which is of interest to the nuclear cluster physics [1]. The best spatial resolution provided by the NTE technique turns out to be a decisive factor for recognition relativistic ⁸Be and ⁹B decays among the projectile fragments [2]. The decays are identified by the invariant mass M^* defined by the sum of all products of 4-momenta P_i of relativistic fragments He and H. Subtracting the sum of the residual masses M is a matter of convenience $Q = M^* - M$. The components P_i are determined by the fragment emission angles under the assumption of conservation a projectile momentum per nucleon. In such an approach the contribution of the decays ⁹B \rightarrow ⁸Be $p \rightarrow 2\alpha p$ in relativistic dissociation of the isotopes ¹⁰B and ^{10,11}C is revealed and, then, an indication to the resonance around 4 MeV in the channel ¹⁰C \rightarrow ⁹Bp found [2].

D. A. Artemenkov · N. K. Kornegrutsa · E. Mitsova · V. V. Rusakova · A. A. Zaitsev · P. I. Zarubin (⊠) · I. G. Zarubina Joint Institute for Nuclear Research (JINR), Dubna, Russia e-mail: zarubin@lhe.jinr.ru

M. Haiduc Institute of Space Science, Magurele, Romania

E. Mitsova · R. Stanoeva Southwestern University, Blagoevgrad, Bulgaria

N. G. Peresadko · A. A. Zaitsev · P. I. Zarubin Lebedev Physical Institute, Russian Academy of Sciences, Moscow, Russia

R. Stanoeva Institute for Nuclear Research and Nuclear Energy, Sofia, Bulgaria

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This experience gave confidence to search in the relativistic dissociation ${}^{12}C \rightarrow 3\alpha$ for the Hoyle state (HS) or the second ${}^{12}C$ exited state exceeding just 378 keV the 3α -threshold [3–5]. The HS studies are reviewed in [6]. This search is inspired by the concept of the α -particle Bose-Einstein condensate whose status is discussed most recently in [6]. In the 90s, the ⁸Be contribution was determined by smallest α -pair opening angles in "white" stars (coherent dissociation not accompanied by any target fragment or produced meson) ${}^{12}C \rightarrow 3\alpha$ [7] and ${}^{16}O \rightarrow 4\alpha$ [8] at energy of 3.65 A GeV. Nowadays, the persisted NTE plates and the relevant data files [7, 8] are reused to obtain distributions over the α -triple invariant mass $Q_{3\alpha}$. Data on the 72 (G.M. Chernov's group, Tashkent) and 114 "white" stars ${}^{12}C \rightarrow 3\alpha$ (A.Sh. Gaitinov's group, Alma-Ata) underwent reanalysis. Besides, additional data on 238 3α stars including 130 "white" ones are obtained recently in this exposure. Then, the NTE layers exposed to 420 A MeV ${}^{12}C$ nuclei allow one to verify the invariant mass approach [5]. In the latter case the α -particle emission angles are measured in the 86 found 3α events including the 36 "white" stars.

The distribution $Q_{3\alpha}$ for all 510 stars is shown in Fig. 24.1. There is a peak in the region $Q_{3\alpha} < 1$ MeV for the 51 stars where HS signal is expected. For events at 3.65 *A* GeV the mean value for the events contributed in the peak $\langle Q_{3\alpha} \rangle$ (RMS) is 397 ± 26 (166) keV, and at 420 *A* MeV, respectively, 346 ± 28 (85) keV. According to the condition $Q_{3\alpha} < 0.7$ MeV 42 (of 424) events at 3.65 *A* GeV can be attributed to HS and 420 *A* MeV–9 (of 86), including 5 "white" stars (of 36). In sum, the contribution of the HS decays to the dissociation of ${}^{12}C \rightarrow 3\alpha$ is $(11 \pm 3)\%$.

HS could emerge in the dissociation ${}^{16}O \rightarrow {}^{12}C^* (\rightarrow 3\alpha) + \alpha$. Figure 24.2 shows the distribution $Q_{3\alpha}$ of all 3α combinations in the 641 "white" stars [8]. While its main part limited $Q_{3\alpha} < 10$ MeV is described by the Rayleigh distribution there is the peak at $Q_{3\alpha} < 700$ keV. The condition $Q_{2\alpha} < 200$ keV meaning at least one ⁸Be per 4 α -event does not affect the statistics in this region. The main part contribution in the peak estimated at 8% is excluded. The remaining 139 events have an average value $\langle Q_{3\alpha} \rangle = (349 \pm 14)$ keV corresponding to HS and RMS 174 keV defined by the

Fig. 24.1 Distribution over invariant mass $Q_{3\alpha}$ of all α -triples in dissociation of ${}^{12}C \rightarrow 3\alpha$ at 3.65 *A* GeV (shaded) and 420 *A* MeV (added by dashed line); line—approximation by the Rayleigh distribution with the parameter $\sigma_{Q_{3\alpha}} = (3.9 \pm 0.4)$ MeV



Fig. 24.2 Distribution over invariant mass $Q_{3\alpha}$ of all 3α combinations in 641 events of "white" stars ¹⁶O \rightarrow 4 α ; line—approximation by the Rayleigh distribution with the parameter $\sigma_{Q_{3\alpha}} = (3.8 \pm 0.2)$ MeV



method resolution. In 9 of them more than one 3α -combination meets the condition $Q_{3\alpha} < 700$ keV. In sum, the contribution of HS decays to coherent dissociation of ${}^{16}\text{O} \rightarrow 4\alpha$ is $(22 \pm 2)\%$. Besides, HS can arise the α -decay product of the ${}^{16}\text{O} 0^+_6$ state [9] (in an analogy with the HS decay ${}^8\text{Be} + \alpha$). The condition $Q_{4\alpha} < 1$ MeV allocates 9 events satisfying with a mean value $\langle Q_{4\alpha} \rangle = (624 \pm 84)$ keV at RMS 252 keV. Then, the estimate of the 0^+_6 contribution is $(7 \pm 2)\%$.

Thus, HS identified in the relativistic ¹²C dissociation is manifested in the ¹⁶O case. These observations indicate that it is not reduced to the unusual ¹²C excitation and, like ⁸Be, is a more universal object of nuclear molecular nature. The closest confirmation of this assumption would be the HS observation in relativistic fragmentation ¹⁴N \rightarrow 3 α . The analysis of the NTE layers exposed to relativistic ¹⁴N nuclei [10] is resumed in the HS context.

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